

EECS 230
ENGINEERING ELECTROMAGNETICS
Leland Pierce

Electrostatics 1

Chapter 4 Overview

Maxwell's Equations

Electrostatics

Magnetostatics

Charge density

Current density

Electric field from charges

Gauss's Law

Electric Scalar Potential Field

Dipole Field

Poisson's eqn

Conductors

current

resistance

joule's law

Dielectrics

polarization

Boundary Conditions

Capacitance

Potential Energy

Image method

$$\mathbf{E} = \int_{v'} d\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{v'} \hat{\mathbf{R}}' \frac{\rho_v dV'}{R'^2}$$

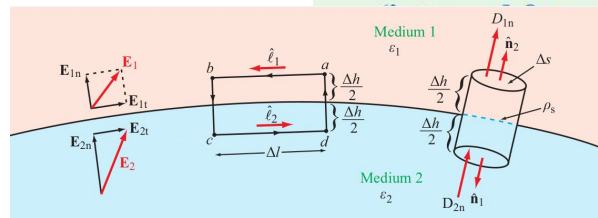
(volume distribution)

$$\nabla \cdot \mathbf{D} = \rho_v,$$

(differential form of Gauss's law)

$$\oint_S \mathbf{D} \cdot d\mathbf{s} = Q$$

(integral form of Gauss's law)



$$V = - \int_{\infty}^P \mathbf{E} \cdot d\mathbf{l}$$

$$\nabla^2 V = - \frac{\rho_v}{\epsilon}$$

$$\mathbf{E} = -\nabla V$$

$$\mathbf{J} = \sigma \mathbf{E} \quad (\text{A/m}^2) \quad \text{(Ohm's law),}$$

$$R = \frac{V}{I} = \frac{- \int_l \mathbf{E} \cdot d\mathbf{l}}{\int_S \mathbf{J} \cdot d\mathbf{s}} = \frac{- \int_l \mathbf{E} \cdot d\mathbf{l}}{\int_S \sigma \mathbf{E} \cdot d\mathbf{s}}$$

$$P = \int_v \mathbf{E} \cdot \mathbf{J} dV \quad (\text{W})$$

$$\nabla \cdot \mathbf{D} = \rho_v,$$

$$\nabla \times \mathbf{E} = - \frac{\partial \mathbf{B}}{\partial t},$$

$$\nabla \cdot \mathbf{B} = 0,$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t}.$$

$$\nabla \cdot \mathbf{D} = \rho_v,$$

$$\nabla \times \mathbf{E} = 0.$$

$$\nabla \cdot \mathbf{B} = 0,$$

$$\nabla \times \mathbf{H} = \mathbf{J}.$$

$$Q = \int_v \rho_v dV \quad (\text{C}).$$

$$\mathbf{J} = \rho_v \mathbf{u} \quad (\text{A/m}^2)$$

$$I = \int_S \mathbf{J} \cdot d\mathbf{s} \quad (\text{A}).$$

Lecture Coverage



Today's lecture:

Sections 4-1 through 4-4 of the book:

4-1: Maxwell's Equations

4-2: Charge and Current Distributions

4-3: Coulomb's Law

4-4: Gauss's Law

4.1 Maxwell's Equations

$$\nabla \cdot \mathbf{D} = \rho_v,$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t},$$

$$\nabla \cdot \mathbf{B} = 0,$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t}.$$

Empirically derived from many measurements

$$\mathbf{D} = \epsilon \mathbf{E}$$

$$\mathbf{B} = \mu \mathbf{H}.$$

E: Electric Field

H: Magnetic Field

J: Current Density

ρ_v : Charge Density

4.1 Maxwell's Equations

$$\nabla \cdot \mathbf{D} = \rho_v,$$

Empirically derived from many measurements

∇

Boldface letters, like **E**, **H**, are vectors and functions of space and time.

∇

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t}.$$

E: Electric Field

H: Magnetic Field

J: Current Density

ρ_v : Charge Density

4.1 Maxwell's Equations

Static Conditions:

Electrostatics

$$\frac{\partial \rho_v}{\partial t} = 0$$

$$\begin{aligned}\nabla \cdot \mathbf{D} &= \rho_v, \\ \nabla \times \mathbf{E} &= 0.\end{aligned}$$

Magnetostatics

$$\frac{\partial \mathbf{J}}{\partial t} = 0$$

$$\begin{aligned}\nabla \cdot \mathbf{B} &= 0, \\ \nabla \times \mathbf{H} &= \mathbf{J}.\end{aligned}$$

Electric and Magnetic Fields are decoupled.

4.2 Charge Distributions

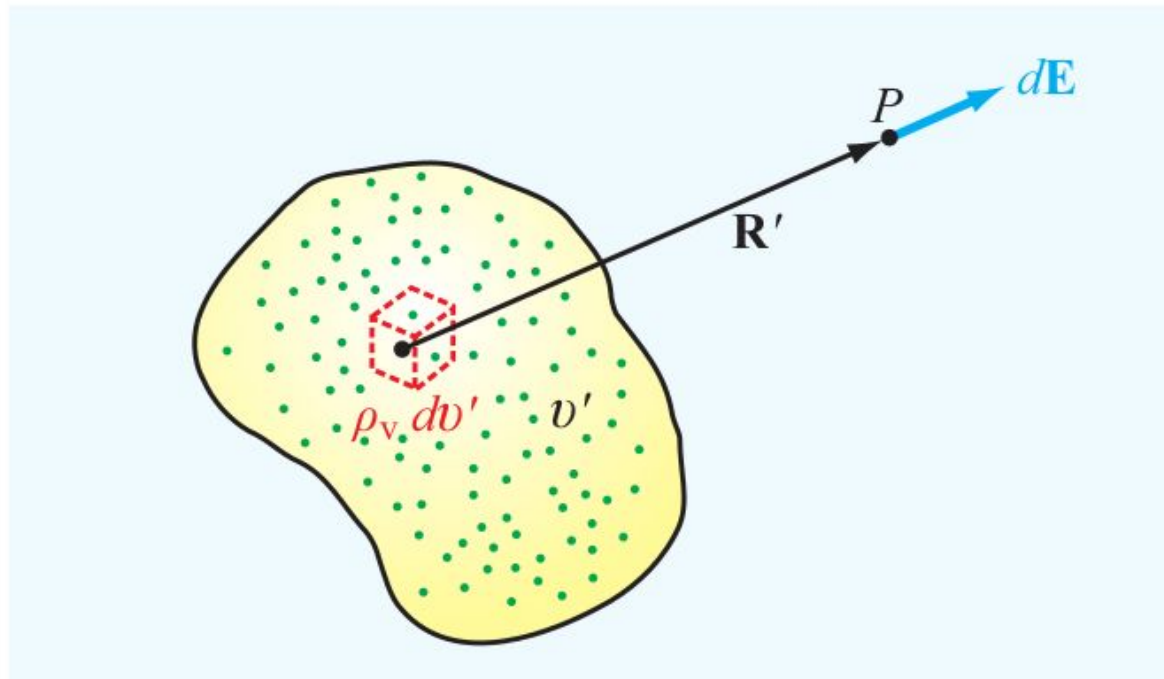
Volume charge density:

$$\rho_v = \lim_{\Delta V \rightarrow 0} \frac{\Delta q}{\Delta V} = \frac{dq}{dV}$$

(C/m³)

Total Charge in a Volume

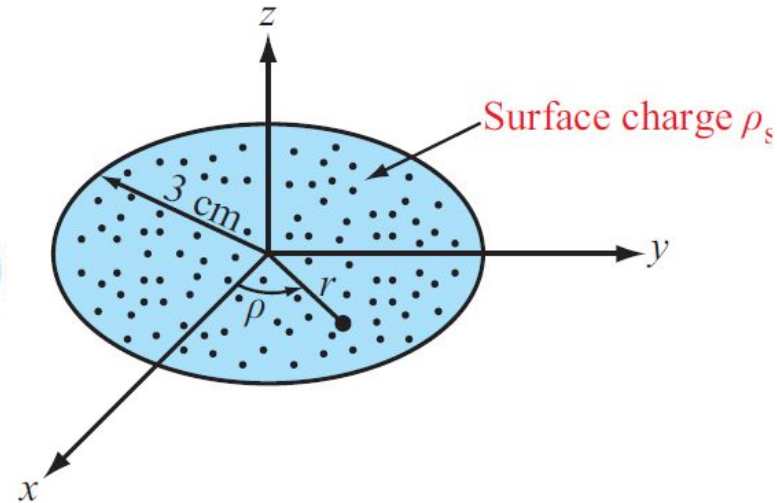
$$Q = \int_V \rho_v dV \quad (C)$$



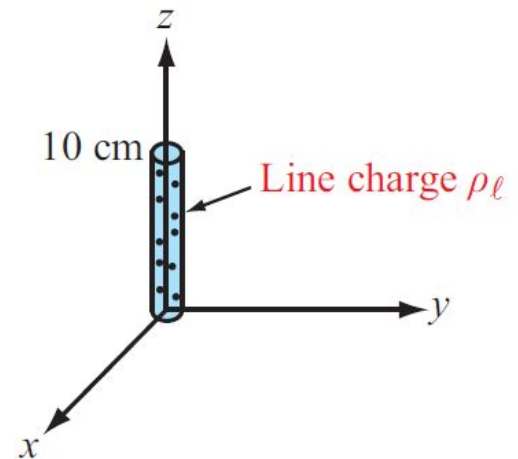
4.2 Charge Distributions

Surface and Line Charge Densities

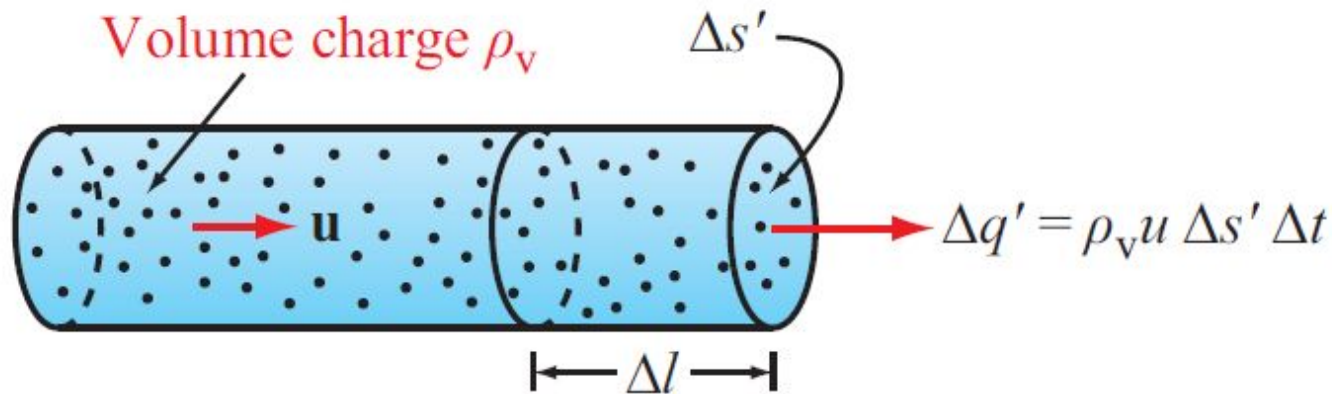
$$\rho_s = \lim_{\Delta s \rightarrow 0} \frac{\Delta q}{\Delta s} = \frac{dq}{ds} \quad (\text{C/m}^2)$$



$$\rho_\ell = \lim_{\Delta l \rightarrow 0} \frac{\Delta q}{\Delta l} = \frac{dq}{dl} \quad (\text{C/m})$$



4.2 Current Density



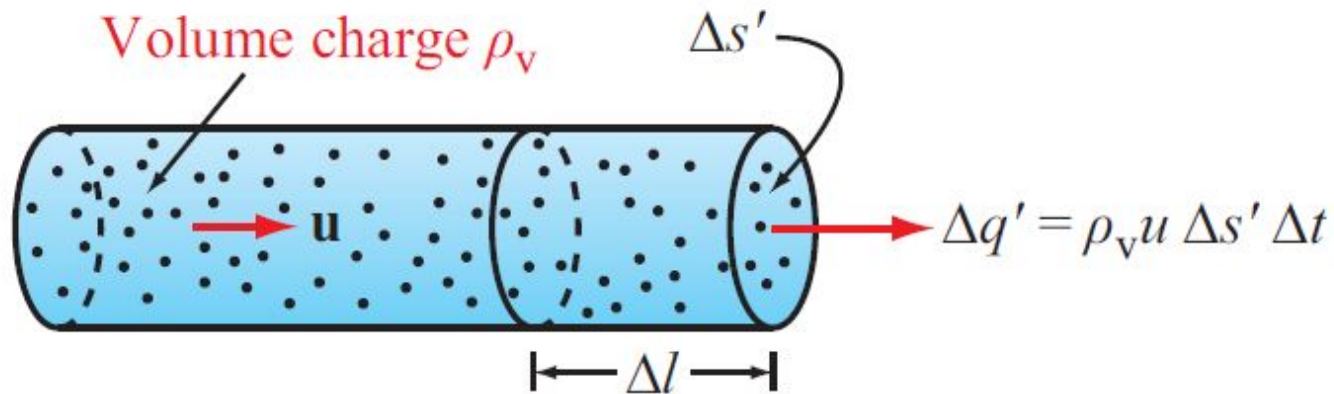
Wire with charge density ρ_v .
cross-sectional area Δs
charges moving with velocity u .

Amount of charge passing a spot
during time interval Δt :

$$\Delta q = \rho_v \mathbf{u} \cdot \Delta \mathbf{s} \Delta t,$$

$$(\text{C/m}^3)(\text{m/sec})(\text{m}^2)(\text{sec})$$

4.2 Current Density



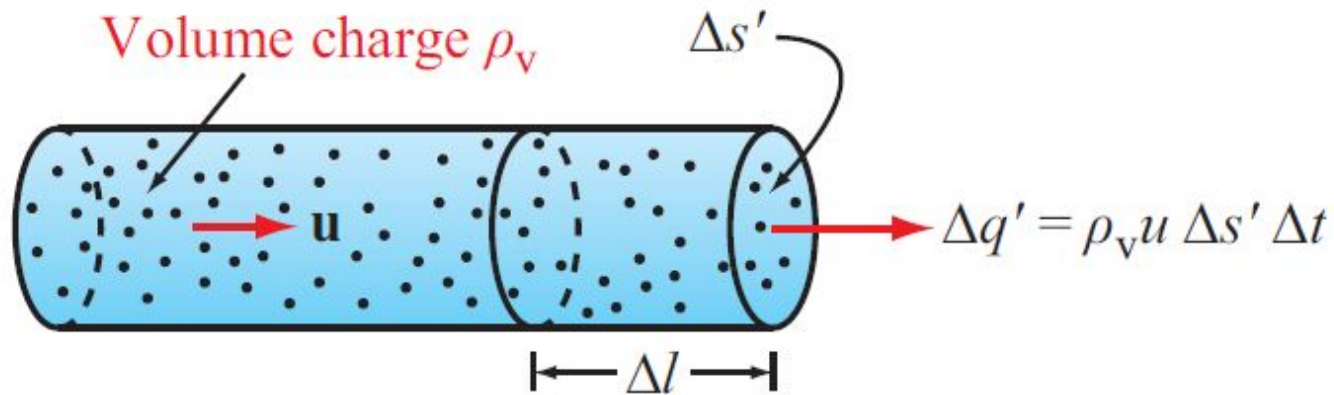
Amount of charge passing a spot during time interval Δt :

$$\Delta q = \rho_v \mathbf{u} \cdot \Delta \mathbf{s} \Delta t.$$

Hence, the current is:

$$\Delta I = \frac{\Delta q}{\Delta t} = \rho_v \mathbf{u} \cdot \Delta \mathbf{s} = \mathbf{J} \cdot \Delta \mathbf{s}.$$

4.2 Current Density



$$\Delta I = \frac{\Delta q}{\Delta t} = \rho_v \mathbf{u} \cdot \Delta \mathbf{s} = \mathbf{J} \cdot \Delta \mathbf{s}.$$

$$\mathbf{J} = \rho_v \mathbf{u} \quad (\text{A/m}^2)$$

\mathbf{J} is called the current density

4.2 Current Density

$$I = \int_S \mathbf{J} \cdot d\mathbf{s} \quad (\text{A}). \quad (4.12)$$

Two kinds of current:

1. **Conduction current:**

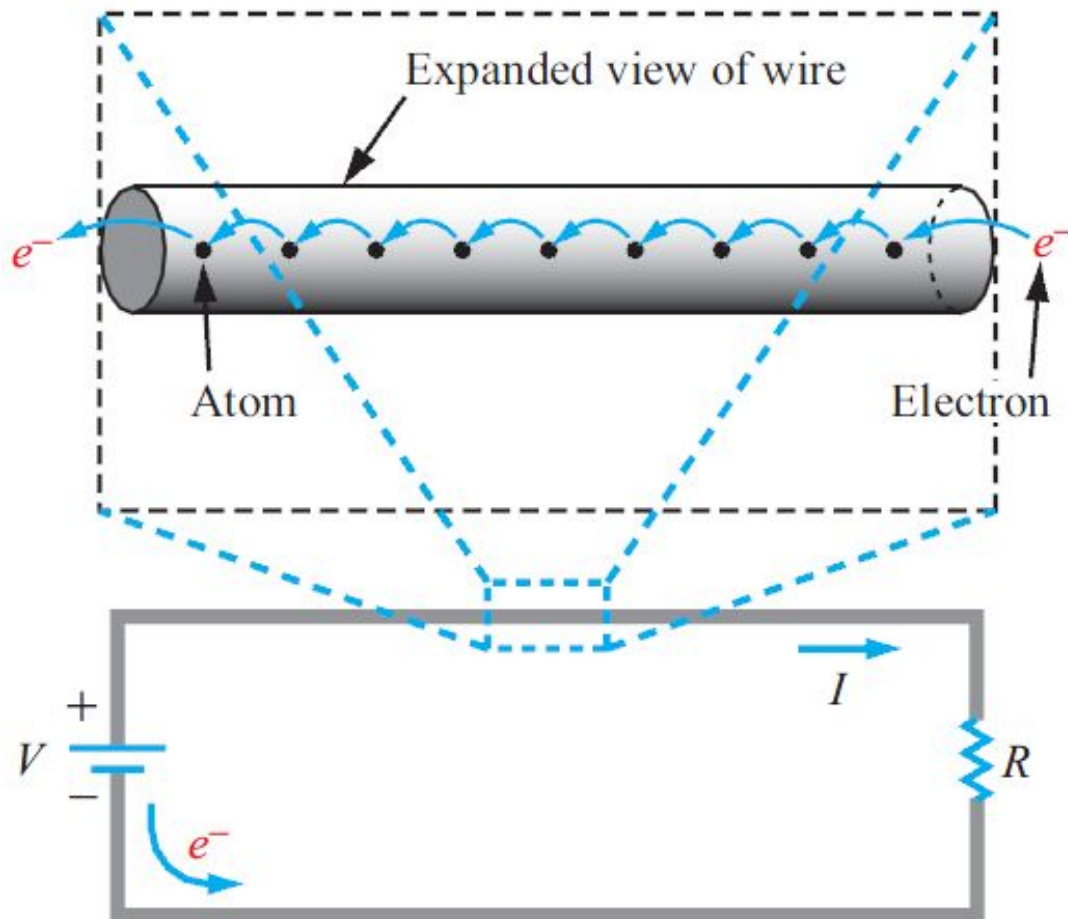
charges moving relative to their host material.
like moving electrons in a wire.

2. **Convection current:**

movement of charged matter

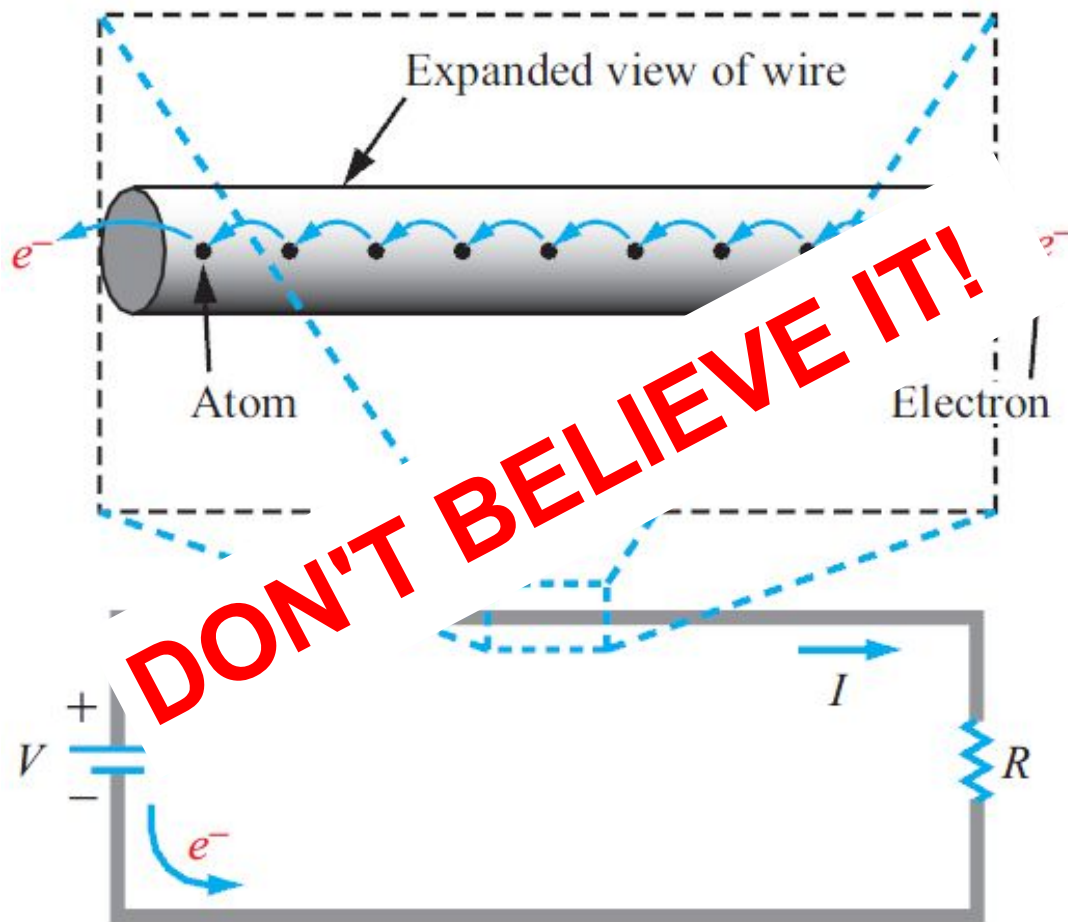
4.2 Convection vs. Conduction

Conduction current: movement of free electrons in a wire



4.2 Convection vs. Conduction

Conduction current: movement of free electrons in a wire



4.2 Quantum Mechanics

Understanding the movement of electrons in matter requires an understanding of quantum mechanics.

This class does not require QM, nor does it use it.

The mathematical models used in this class are "classical", applying only to large-scale situations, and hence incorrect or misleading when it comes to the very small.

Use **Quantum Electromagnetics** concepts and tools when your application involves situations that are very small.

4.3 Coulomb's Law

Electric field at point P due to single charge

$$\mathbf{E} = \hat{\mathbf{R}} \frac{q}{4\pi\epsilon R^2} \quad (\text{V/m})$$

Electric force on a test charge placed at P

$$\mathbf{F} = q'\mathbf{E} \quad (\text{N})$$

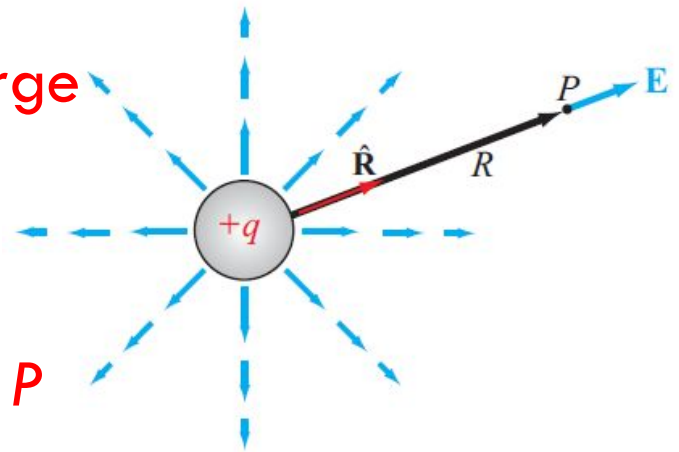


Figure 4-3: Electric-field lines due to a charge q .

4.3 Coulomb's Law

Electric flux density **D**

$$\mathbf{D} = \epsilon \mathbf{E}$$

$$\epsilon = \epsilon_r \epsilon_0$$

$$\epsilon_0 = 8.85 \times 10^{-12} \approx (1/36\pi) \times 10^{-9} \quad (\text{F/m})$$

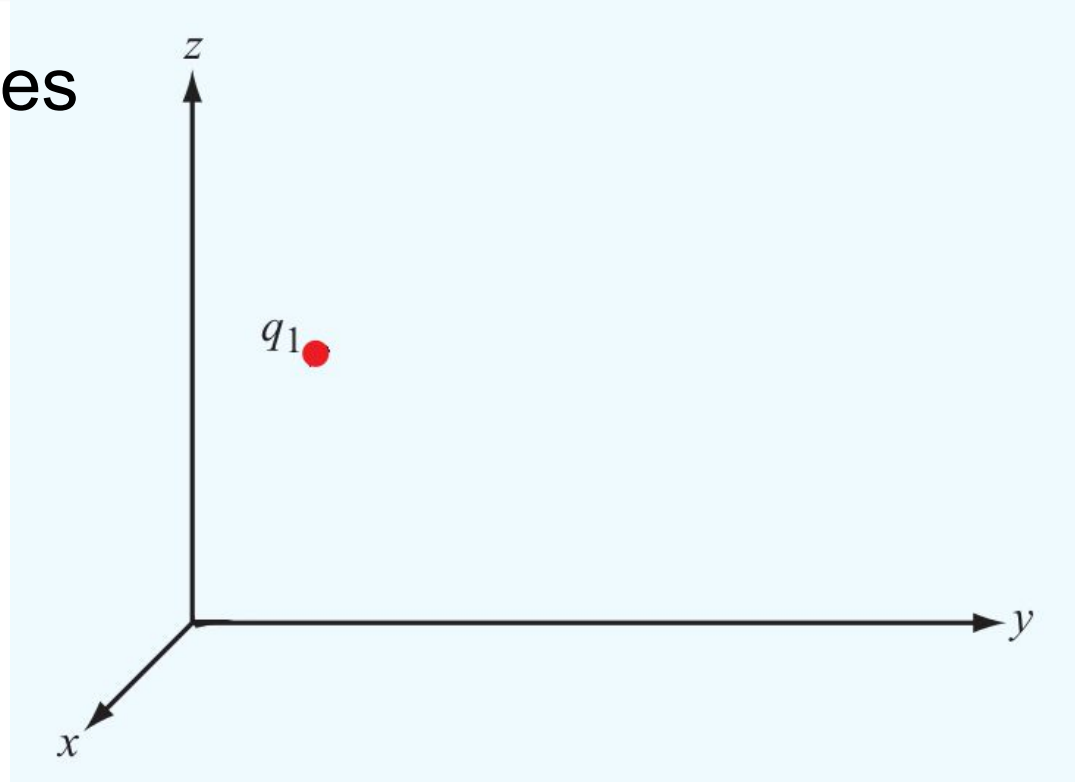
Material is:

Linear: $\epsilon \neq f(|\mathbf{E}|)$

Isotropic: $\epsilon \neq f(\text{direction of } \mathbf{E})$

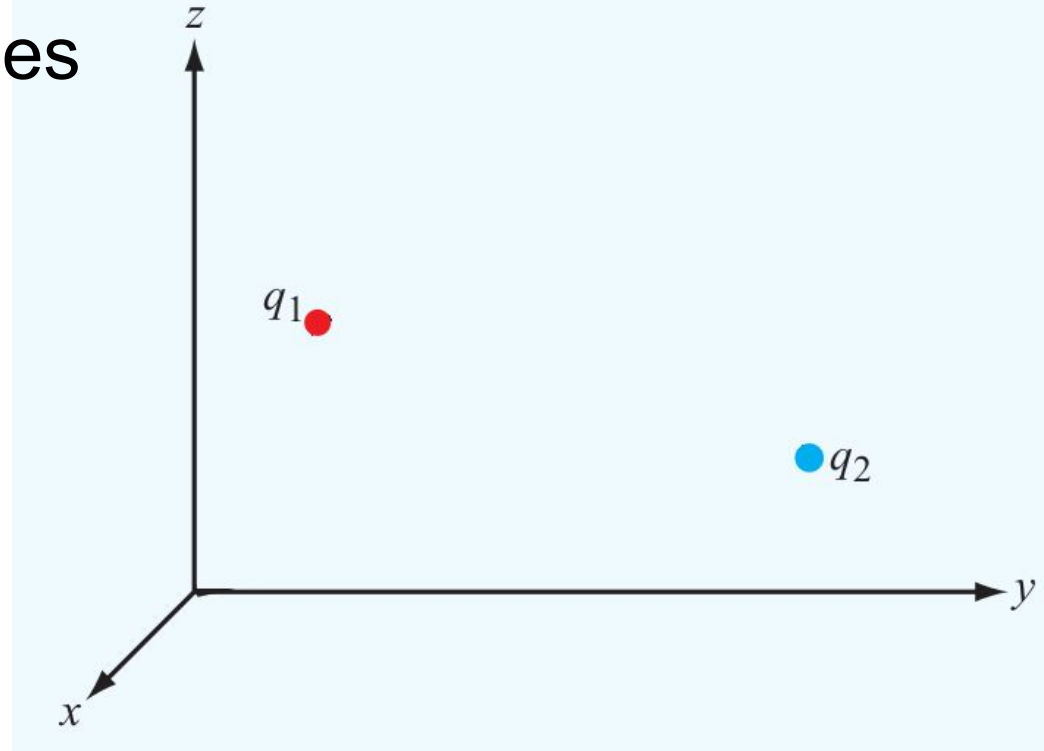
4.3 Electric Field of Multiple Charges

Start by placing 2 charges somewhere in space:



4.3 Electric Field of Multiple Charges

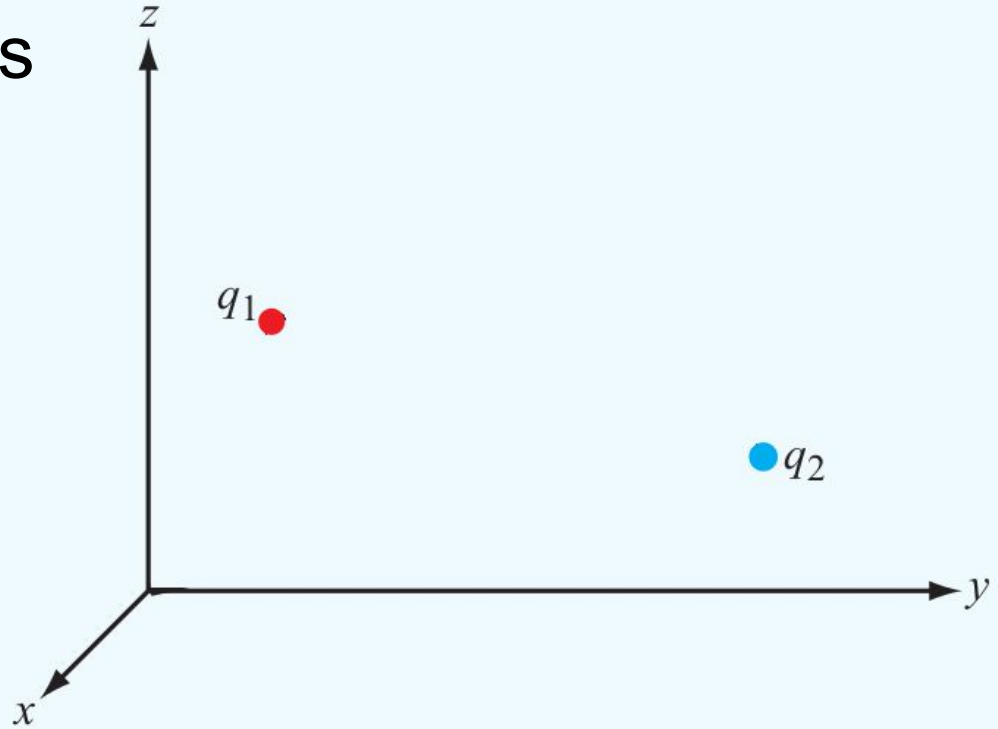
Start by placing 2 charges somewhere in space:



4.3 Electric Field of Multiple Charges

Start by placing 2 charges somewhere in space:

Next, identify a point P in space where we want to determine the \mathbf{E} field.

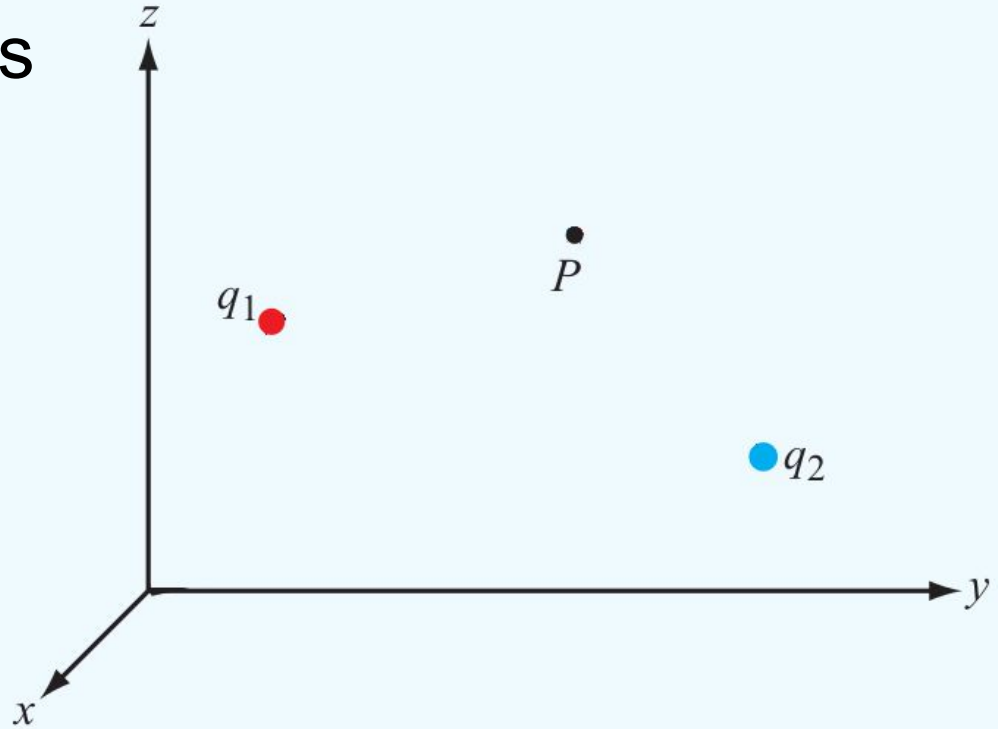


4.3 Electric Field of Multiple Charges

Start by placing 2 charges somewhere in space:

Next, identify a point P in space where we want to determine the \mathbf{E} field.

The position of point P is given by a vector \mathbf{R} from the origin

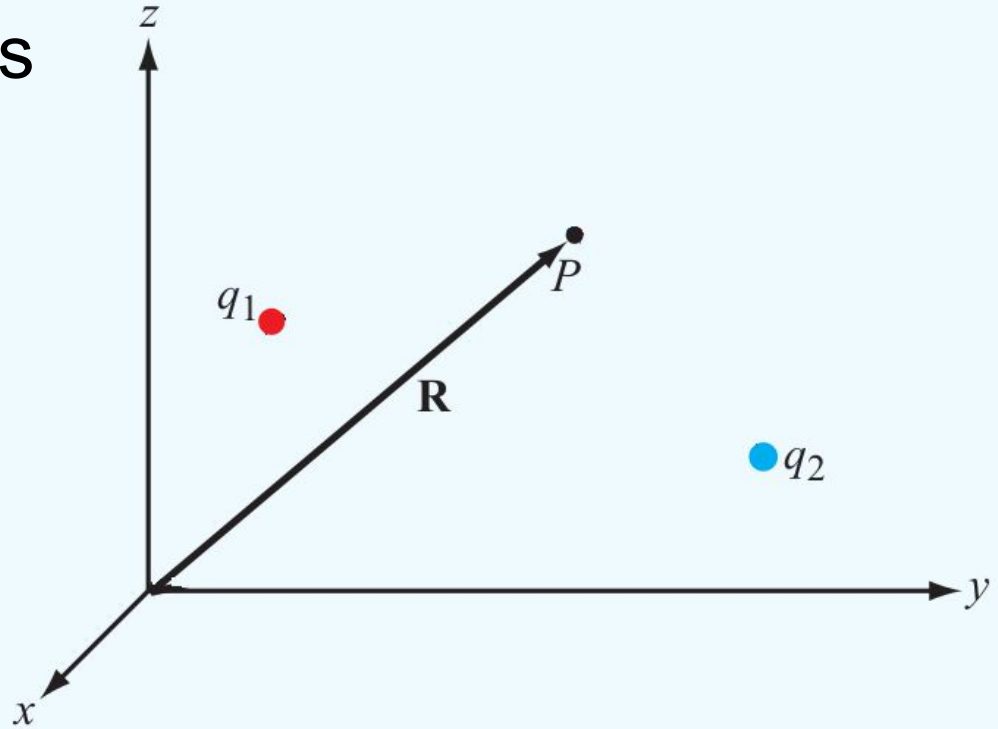


4.3 Electric Field of Multiple Charges

Start by placing 2 charges somewhere in space:

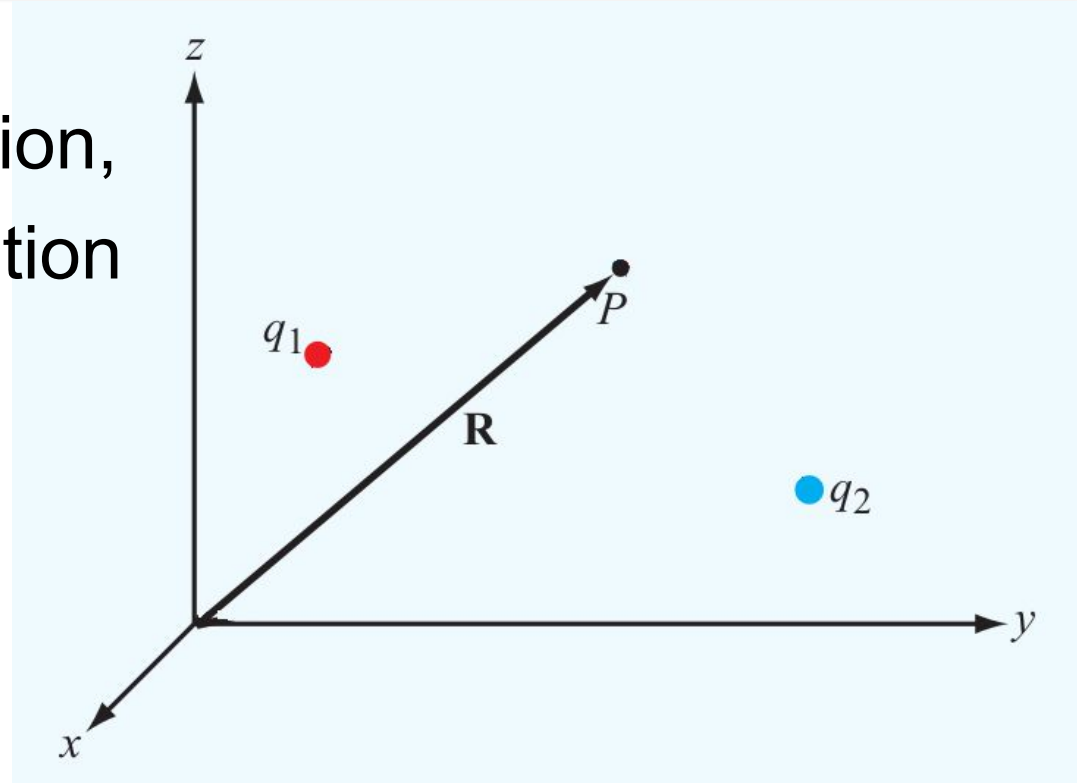
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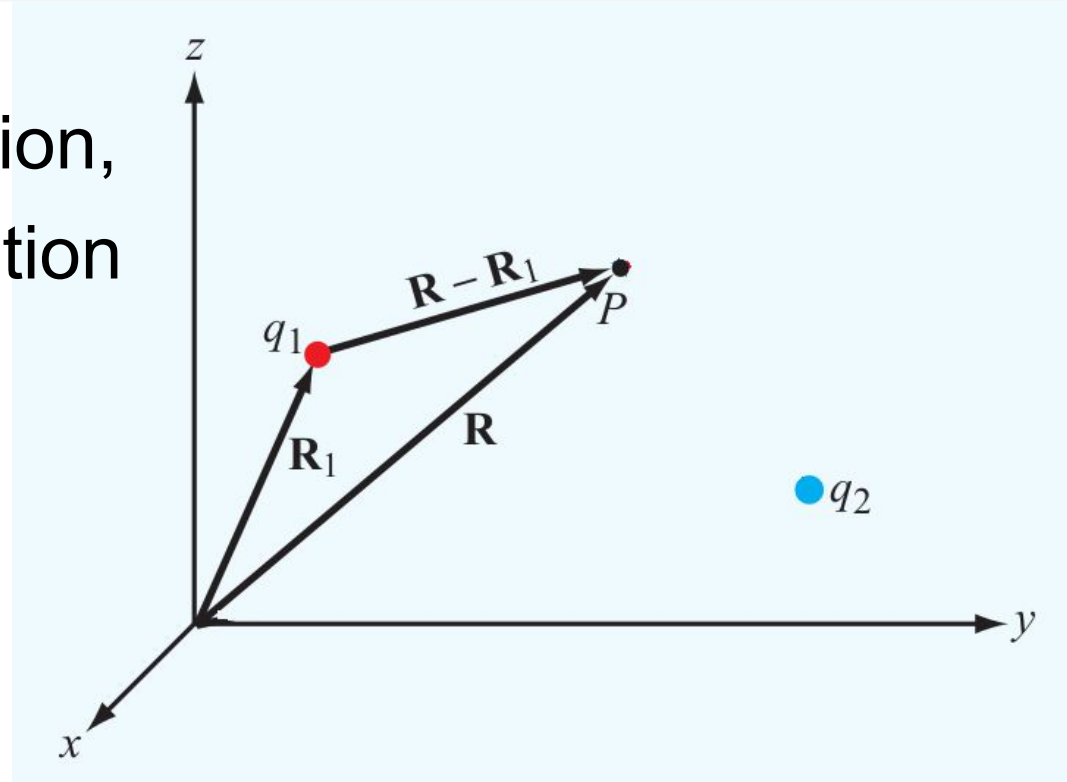
4.3 Electric Field of Multiple Charges

To get \mathbf{E} due to q_1 we need to identify its position, \mathbf{R}_1 , and its *relative* position to point P : $\mathbf{R}-\mathbf{R}_1$



4.3 Electric Field of Multiple Charges

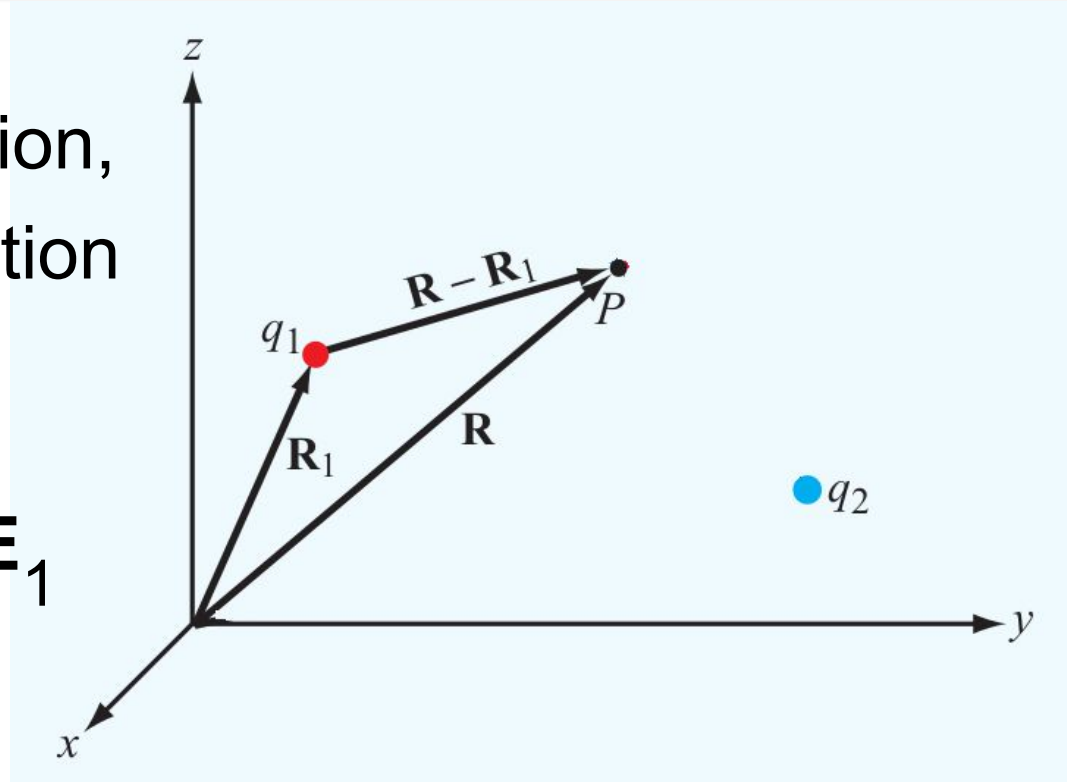
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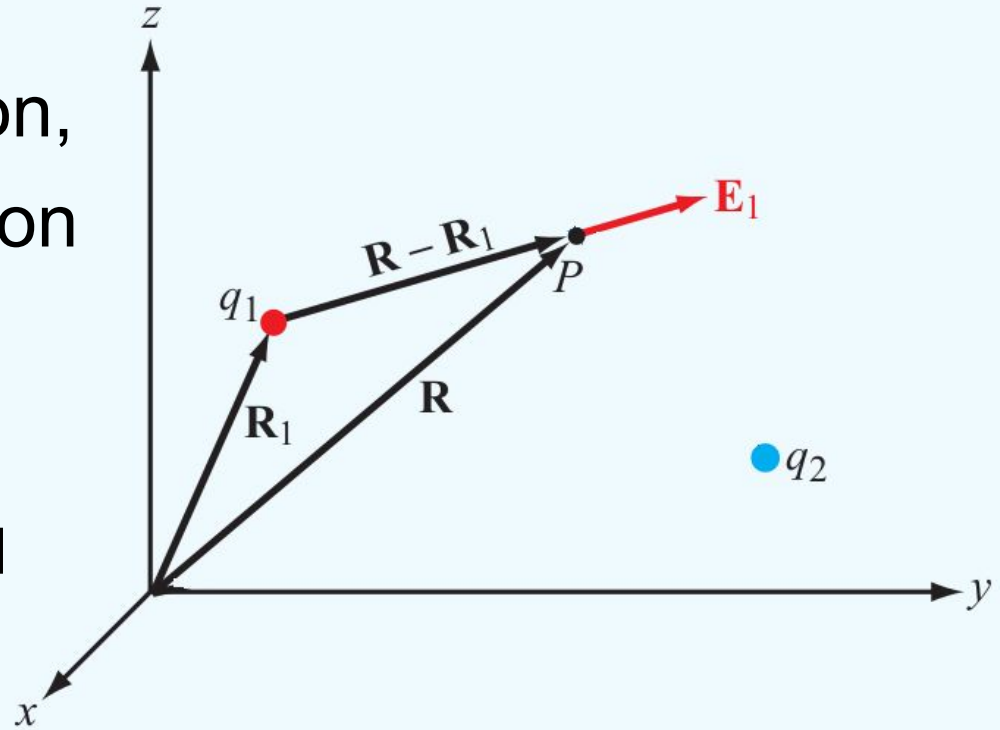
Calculate \mathbf{E} due to q_1 : \mathbf{E}_1



4.3 Electric Field of Multiple Charges

To get \mathbf{E} due to q_1 we need to identify its position, \mathbf{R}_1 , and its *relative* position to point P : $\mathbf{R} - \mathbf{R}_1$

Calculate \mathbf{E} due to q_1 : \mathbf{E}_1

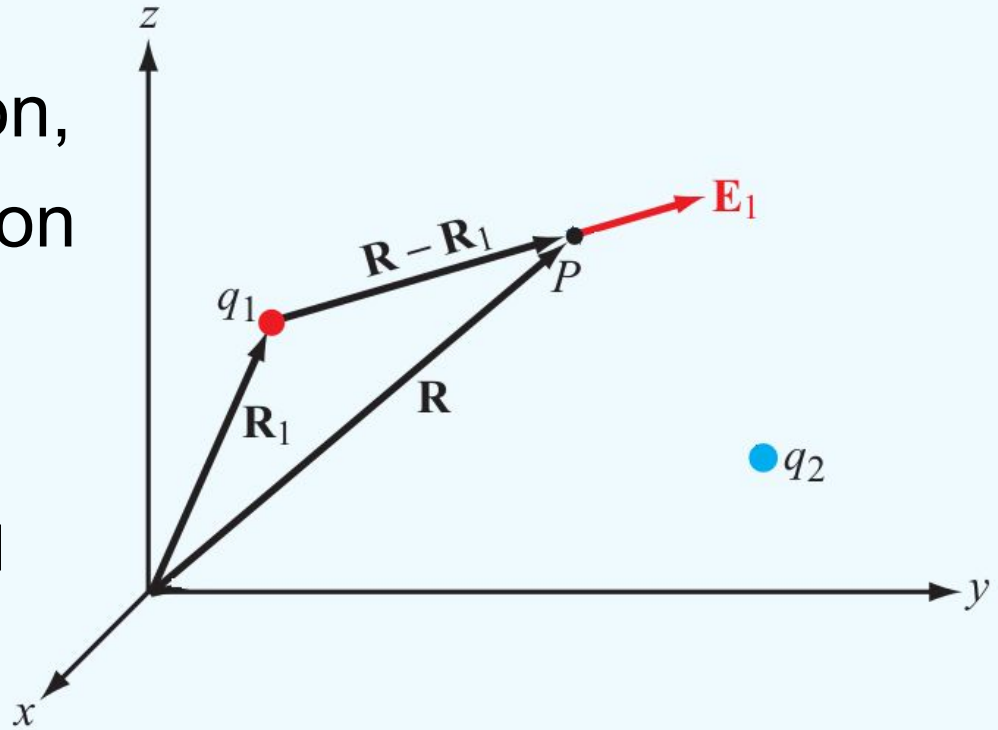


4.3 Electric Field of Multiple Charges

To get \mathbf{E} due to q_1 we need to identify its position, \mathbf{R}_1 , and its *relative* position to point P : $\mathbf{R}-\mathbf{R}_1$

Calculate \mathbf{E} due to q_1 : \mathbf{E}_1

To get \mathbf{E} due to q_2 we need to identify its position, \mathbf{R}_2 , and its *relative* position to point P : $\mathbf{R}-\mathbf{R}_2$

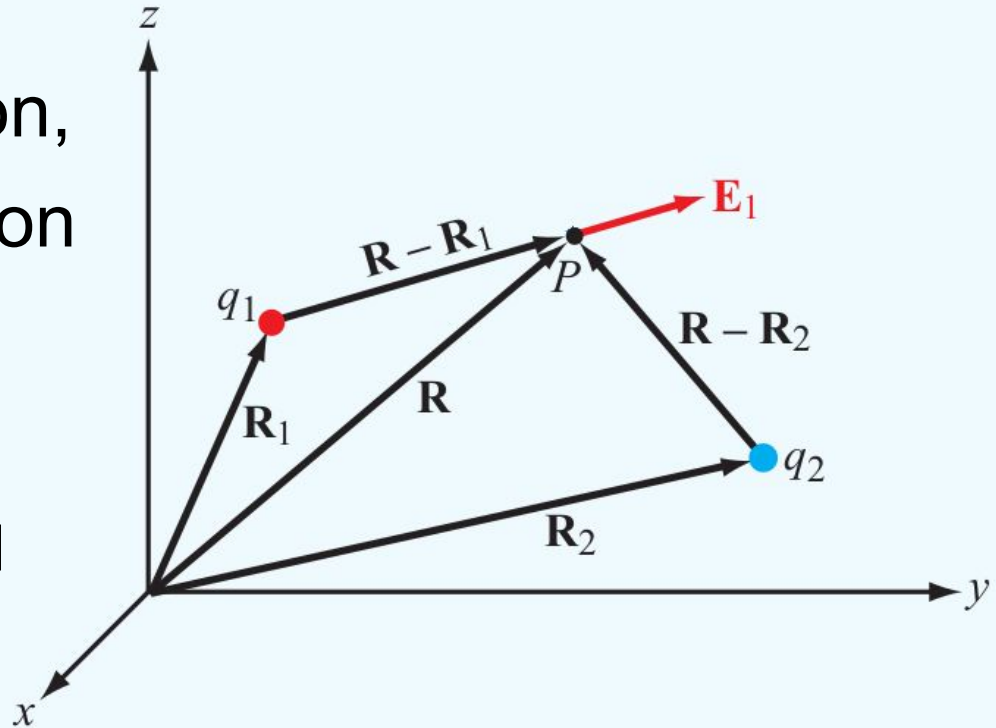


4.3 Electric Field of Multiple Charges

To get \mathbf{E} due to q_1 we need to identify its position, \mathbf{R}_1 , and its *relative* position to point P : $\mathbf{R}-\mathbf{R}_1$

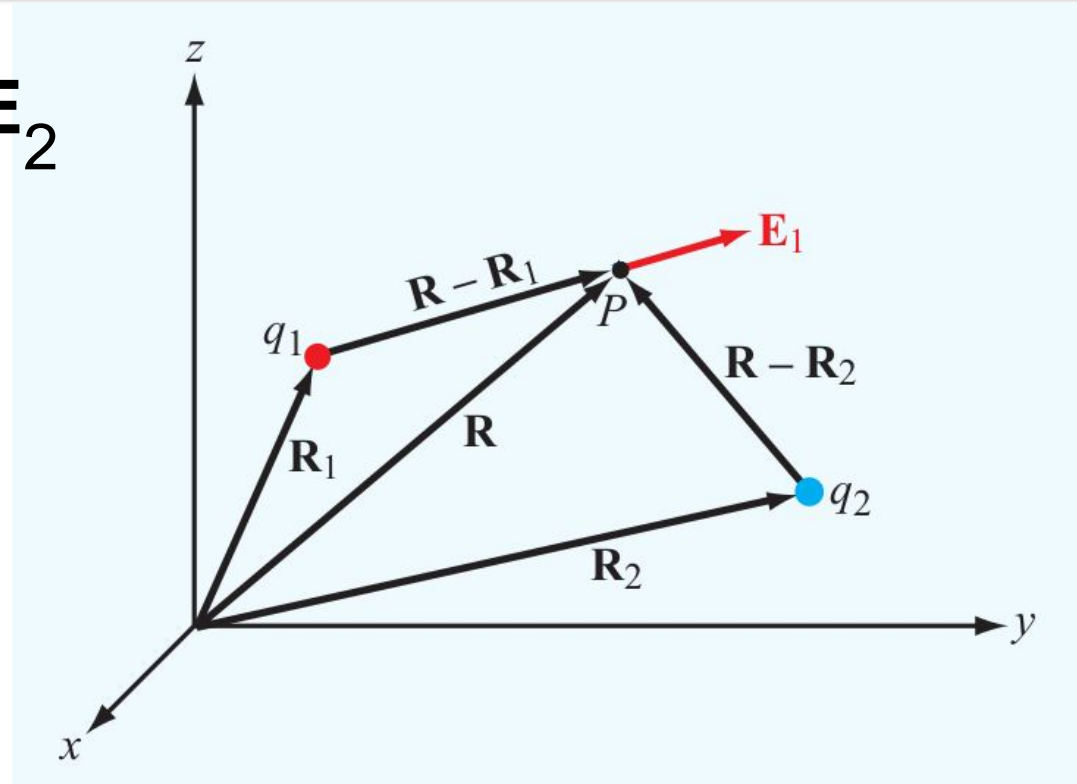
Calculate \mathbf{E} due to q_1 : \mathbf{E}_1

To get \mathbf{E} due to q_2 we need to identify its position, \mathbf{R}_2 , and its *relative* position to point P : $\mathbf{R}-\mathbf{R}_2$



4.3 Electric Field of Multiple Charges

Calculate \mathbf{E} due to q_2 : \mathbf{E}_2

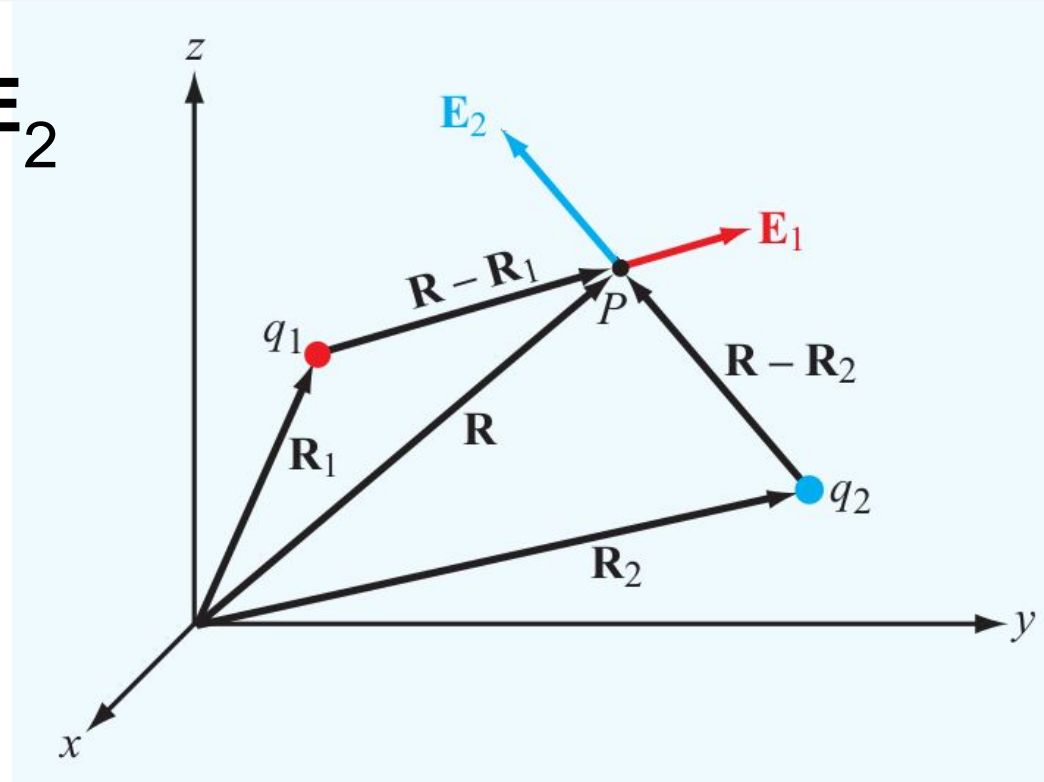


4.3 Electric Field of Multiple Charges

Calculate \mathbf{E} due to q_2 : \mathbf{E}_2

Total \mathbf{E} is sum:

$$\mathbf{E} = \mathbf{E}_1 + \mathbf{E}_2$$

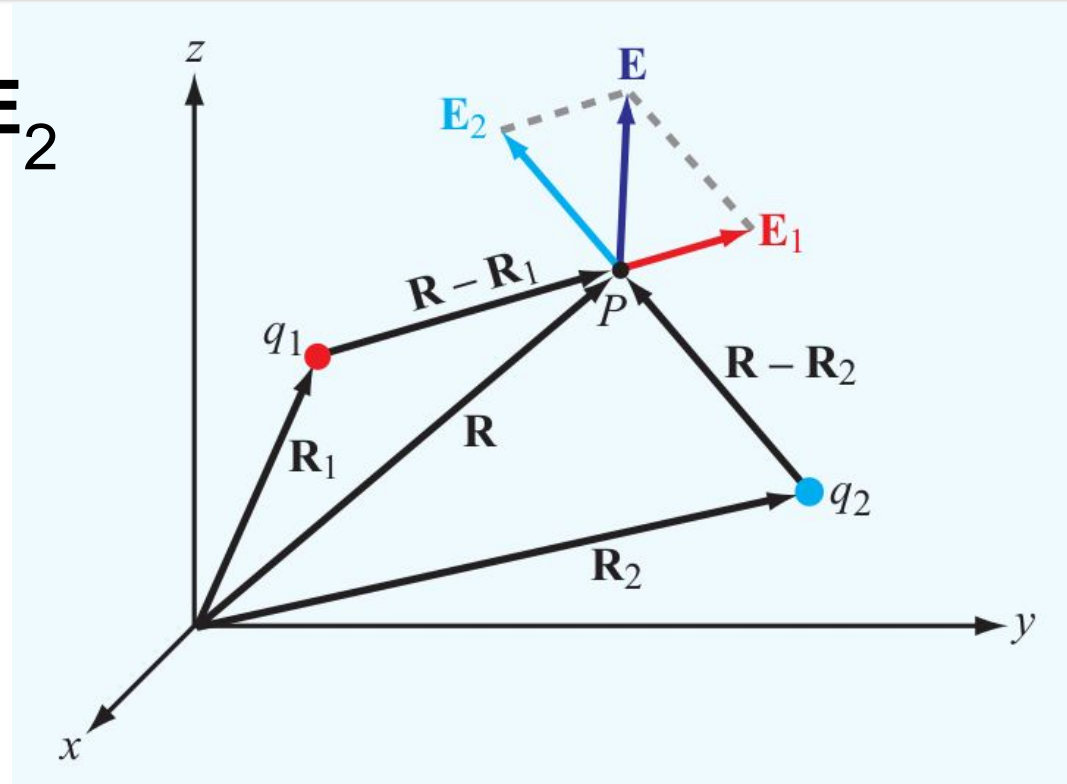


4.3 Electric Field of Multiple Charges

Calculate \mathbf{E} due to q_2 : \mathbf{E}_2

Total \mathbf{E} is sum:

$$\mathbf{E} = \mathbf{E}_1 + \mathbf{E}_2$$



4.3 Electric Field of Multiple Charges

In general:

$$\mathbf{E} = \hat{\mathbf{R}} \frac{q}{4\pi\epsilon R^2} \quad (\text{V/m})$$

SO:

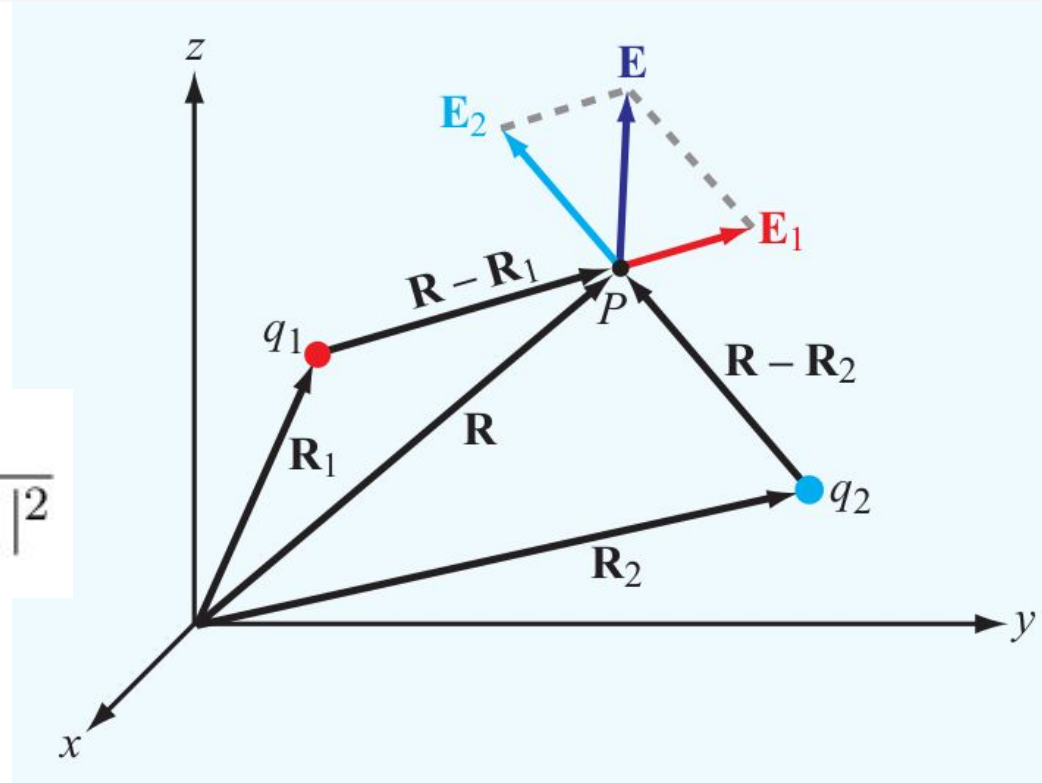
$$\mathbf{E}_1 = \widehat{\mathbf{R} - \mathbf{R}_1} \frac{q_1}{4\pi\epsilon |\mathbf{R} - \mathbf{R}_1|^2}$$

where:

$$\widehat{\mathbf{R} - \mathbf{R}_1} = \frac{\mathbf{R} - \mathbf{R}_1}{|\mathbf{R} - \mathbf{R}_1|}$$

SO:

$$\mathbf{E}_1 = (\mathbf{R} - \mathbf{R}_1) \frac{q_1}{4\pi\epsilon |\mathbf{R} - \mathbf{R}_1|^3}$$



4.3 Electric Field of Multiple Charges

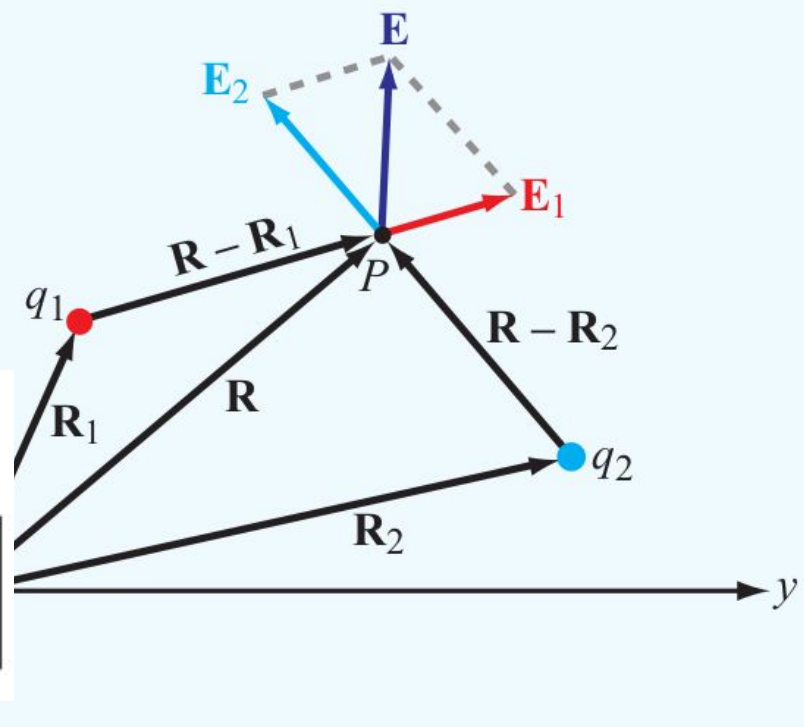
Similarly:

$$\mathbf{E}_2 = (\mathbf{R} - \mathbf{R}_2) \frac{q_2}{4\pi\epsilon |\mathbf{R} - \mathbf{R}_2|^3}$$

Hence:

$$\mathbf{E} = \mathbf{E}_1 + \mathbf{E}_2$$

$$= \frac{1}{4\pi\epsilon} \left[\frac{q_1(\mathbf{R} - \mathbf{R}_1)}{|\mathbf{R} - \mathbf{R}_1|^3} + \frac{q_2(\mathbf{R} - \mathbf{R}_2)}{|\mathbf{R} - \mathbf{R}_2|^3} \right]$$



For N charges:

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \sum_{i=1}^N \frac{q_i(\mathbf{R} - \mathbf{R}_i)}{|\mathbf{R} - \mathbf{R}_i|^3}$$

Ex. 4-3 Electric Field of 2 Charges

Given: in free space:

$$q_1 = 2 \times 10^{-5} \text{ C}$$

$$q_2 = -4 \times 10^{-5} \text{ C}$$

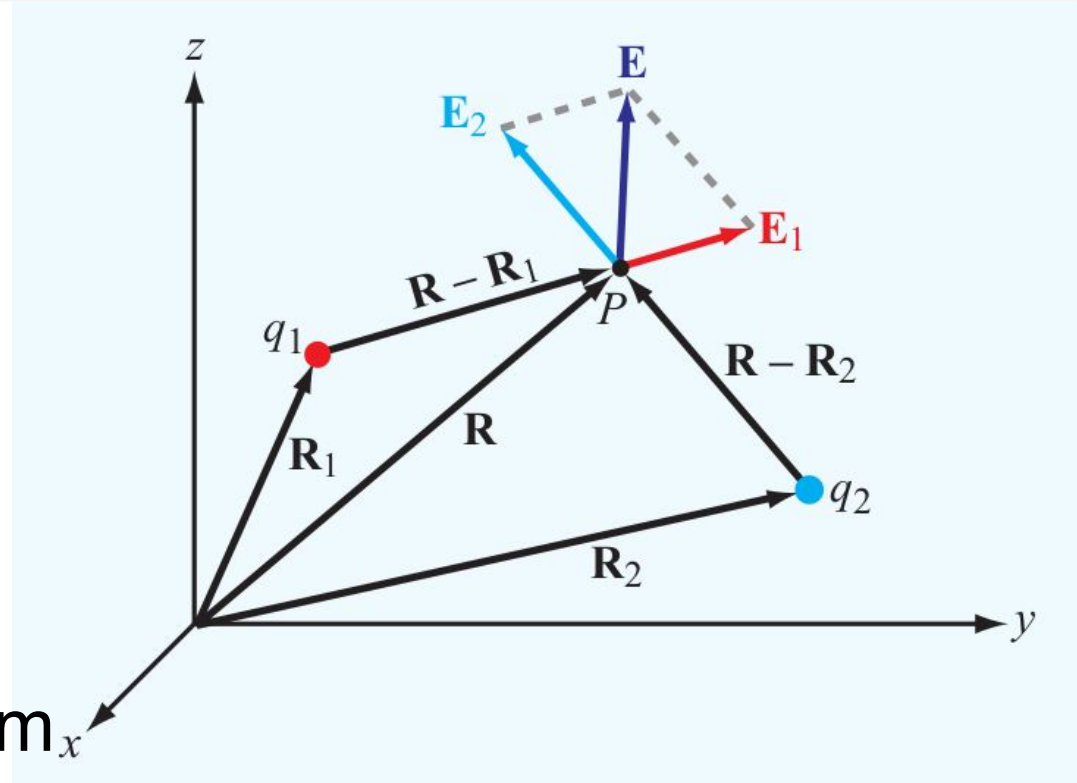
$$P_1 = (1, 3, -1) \text{ meters}$$

$$P_2 = (-3, 1, -2) \text{ meters}$$

Find: \mathbf{E} at $P = (3, 1, -2) \text{ m}_x$

Solution:

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \sum_{i=1}^N \frac{q_i(\mathbf{R} - \mathbf{R}_i)}{|\mathbf{R} - \mathbf{R}_i|^3}$$



Ex. 4-3 Electric Field of 2 Charges

Solution:

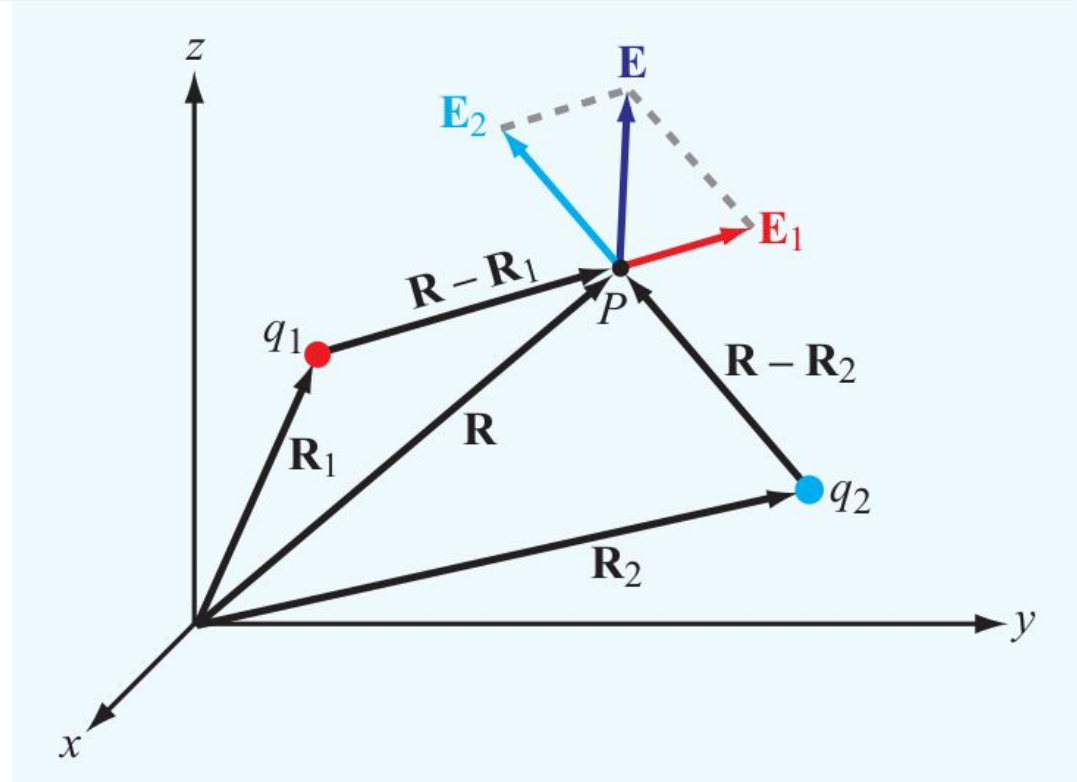
$$\mathbf{E} = \frac{1}{4\pi\epsilon} \sum_{i=1}^N \frac{q_i(\mathbf{R} - \mathbf{R}_i)}{|\mathbf{R} - \mathbf{R}_i|^3}$$

so need:

$$\mathbf{R}_1 = \hat{x}3 + \hat{y}3 - \hat{z},$$

$$\mathbf{R}_2 = -\hat{x}3 + \hat{y} - \hat{z}2,$$

$$\mathbf{R} = \hat{x}3 + \hat{y} - \hat{z}2.$$



Ex. 4-3 Electric Field of 2 Charges

Solution:

$$\mathbf{R}_1 = \hat{x}3 + \hat{y}3 - \hat{z},$$

$$\mathbf{R}_2 = -\hat{x}3 + \hat{y} - \hat{z}2,$$

$$\mathbf{R} = \hat{x}3 + \hat{y} - \hat{z}2.$$

$$\mathbf{R} - \mathbf{R}_1 = \hat{x}3 + \hat{y} - \hat{z}2 - (\hat{x}3 + \hat{y}3 - \hat{z})$$

$$\mathbf{R} - \mathbf{R}_1 = \hat{x}2 - \hat{y}2 - \hat{z}$$

$$|\mathbf{R} - \mathbf{R}_1| = \sqrt{2^2 + 2^2 + 1^2} = \sqrt{9} = 3$$

$$|\mathbf{R} - \mathbf{R}_1|^3 = 3^3 = 27$$

Ex. 4-3 Electric Field of 2 Charges

Solution:

$$\mathbf{R}_1 = \hat{x}3 + \hat{y}3 - \hat{z},$$

$$\mathbf{R}_2 = -\hat{x}3 + \hat{y} - \hat{z}2,$$

$$\mathbf{R} = \hat{x}3 + \hat{y} - \hat{z}2.$$

$$\mathbf{R} - \mathbf{R}_2 = \hat{x}3 + \hat{y} - \hat{z}2 - (-\hat{x}3 + \hat{y} - \hat{z}2)$$

$$\mathbf{R} - \mathbf{R}_2 = \hat{x}6 + \hat{y}0 - \hat{z}0$$

$$\mathbf{R} - \mathbf{R}_2 = \hat{x}6$$

$$|\mathbf{R} - \mathbf{R}_2| = 6$$

$$|\mathbf{R} - \mathbf{R}_2|^3 = 216$$

Ex. 4-3 Electric Field of 2 Charges

Solution: plug in:

$$\mathbf{E} = \frac{1 \times 10^{-5} \text{C}}{4\pi\epsilon_0} \left[\frac{2(\hat{x}^2 - \hat{y}^2 - \hat{z})}{27} - \frac{4\hat{x}6}{216} \right]$$

$$\mathbf{E} = \frac{1 \times 10^{-5} \text{C}}{4\pi\epsilon_0} \left[\frac{2(216)(\hat{x}^2 - \hat{y}^2 - \hat{z})}{27(216)} - \frac{4(27)\hat{x}6}{27(216)} \right]$$

$$\mathbf{E} = \frac{1 \times 10^{-5} \text{C}}{4\pi\epsilon_0} \left[\frac{\hat{x}864 - \hat{y}864 - \hat{z}432}{5832} - \frac{\hat{x}648}{5832} \right]$$

$$\mathbf{E} = \frac{1 \times 10^{-5} \text{C}}{4\pi\epsilon_0} \left[\frac{\hat{x}1 - \hat{y}4 - \hat{z}2}{27} \right]$$

Ex. 4-3 Electric Field of 2 Charges

Solution: plug in:

$$\mathbf{E} = \frac{1 \times 10^{-5} \text{C}}{4\pi\epsilon_0} \left[\frac{\hat{\mathbf{x}}1 - \hat{\mathbf{y}}4 - \hat{\mathbf{z}}2}{27} \right]$$

$$\mathbf{E} = \frac{1 \times 10^{-5} \text{C}}{108\pi(8.85 \times 10^{-12} \text{F/m})} (\hat{\mathbf{x}}1 - \hat{\mathbf{y}}4 - \hat{\mathbf{z}}2)$$

$$\mathbf{E} = 3330(\hat{\mathbf{x}} - \hat{\mathbf{y}}4 - \hat{\mathbf{z}}2) \text{V/m}$$

Ex. 4-3 Electric Field of 2 Charges

Solution: to calculate the force on $q_3 = 8 \times 10^{-5} \text{ C}$

$$\mathbf{F} = q\mathbf{E}$$

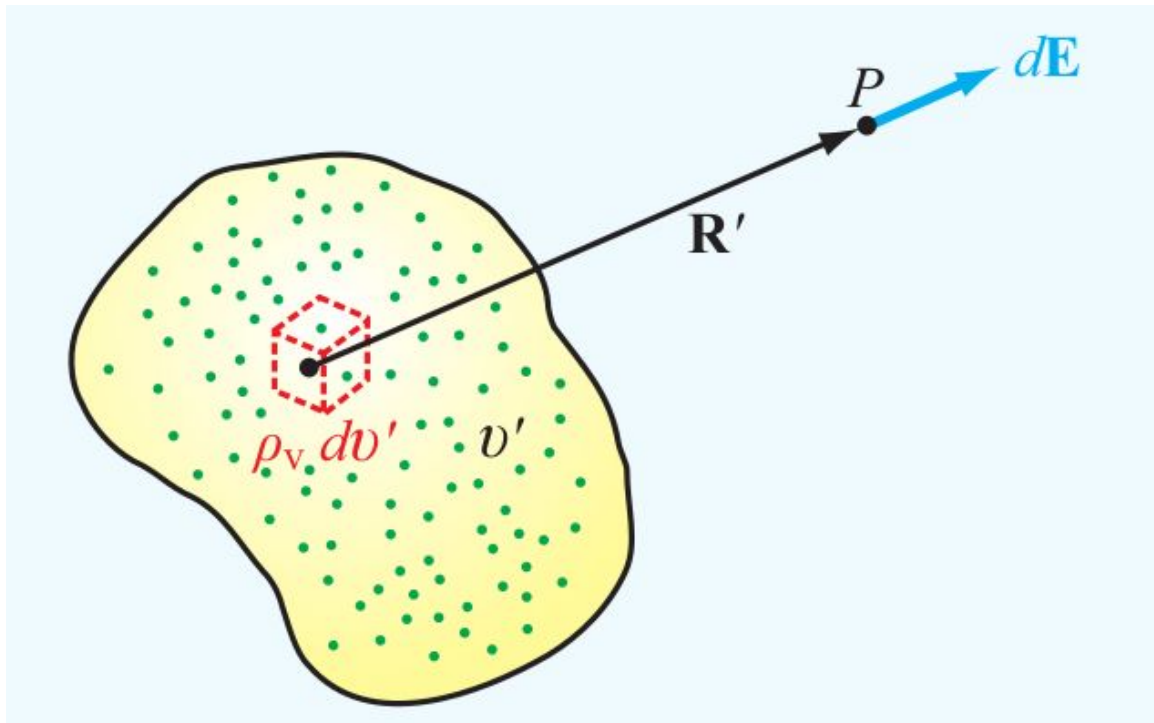
$$\mathbf{F} = q_3\mathbf{E}$$

$$\mathbf{F} = (8 \times 10^{-5} \text{ C}) 3330(\hat{\mathbf{x}} - \hat{\mathbf{y}}4 - \hat{\mathbf{z}}2) \text{ V/m}$$

$$\mathbf{F} = 266.4(\hat{\mathbf{x}} - \hat{\mathbf{y}}4 - \hat{\mathbf{z}}2) \text{ mN}$$

4.3 Electric Field Due to Charge Distributions

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{v'} \hat{\mathbf{R}}' \frac{\rho_v dV'}{R'^2} \quad (\text{volume distribution})$$



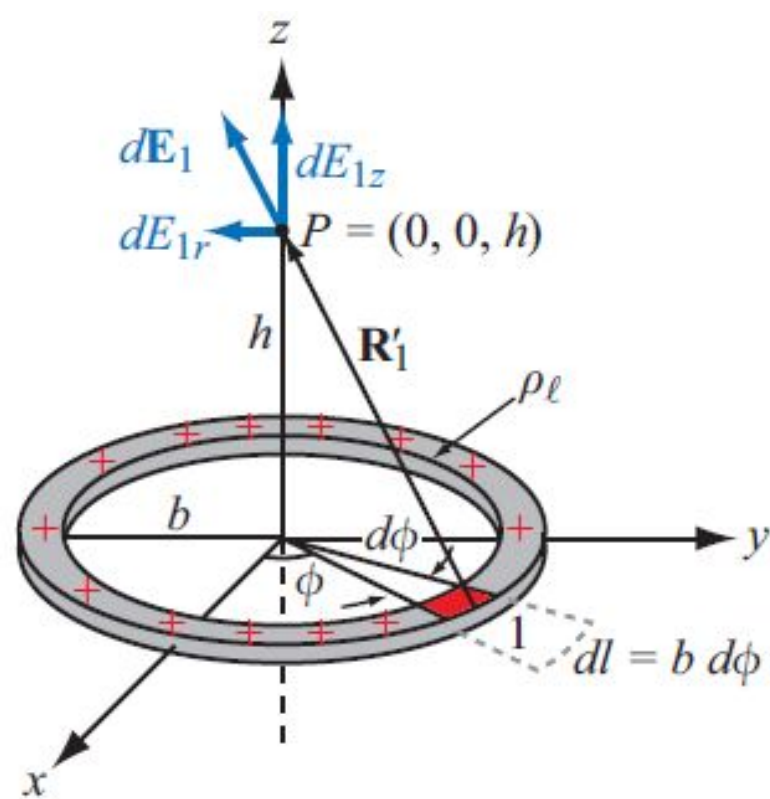
4.3 Electric Field Due to Charge Distributions

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{\mathcal{V}'} \hat{\mathbf{R}}' \frac{\rho_v d\mathcal{V}'}{R'^2} \quad \text{(volume distribution)}$$

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{S'} \hat{\mathbf{R}}' \frac{\rho_s ds'}{R'^2} \quad \text{(surface distribution)}$$

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{l'} \hat{\mathbf{R}}' \frac{\rho_l dl'}{R'^2} \quad \text{(line distribution)}$$

Example 4-4

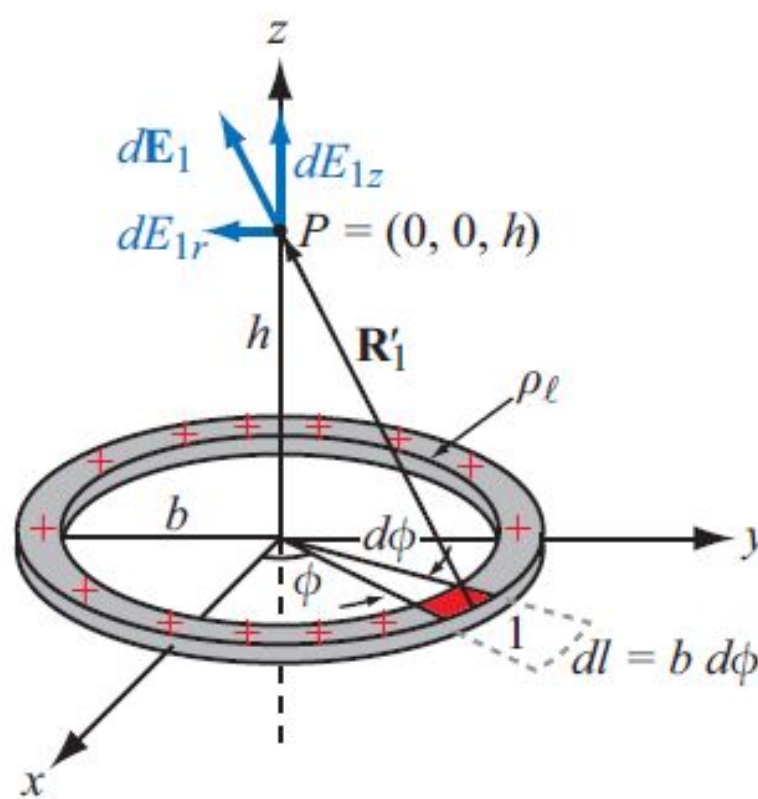


Given: thin ring: radius = b
line charge density = ρ_l
in x - y plane,

Find: $\mathbf{E}(0,0,h)$
(along the axis at distance h)

Example 4-4

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{l'} \hat{\mathbf{R}}' \frac{\rho_{\ell} dl'}{R'^2}$$



Solution:

Find \mathbf{E} due to small segment of charge:

center: $(b, \phi, 0)$ (Cylindrical coords)

length: $dl = b d\phi$

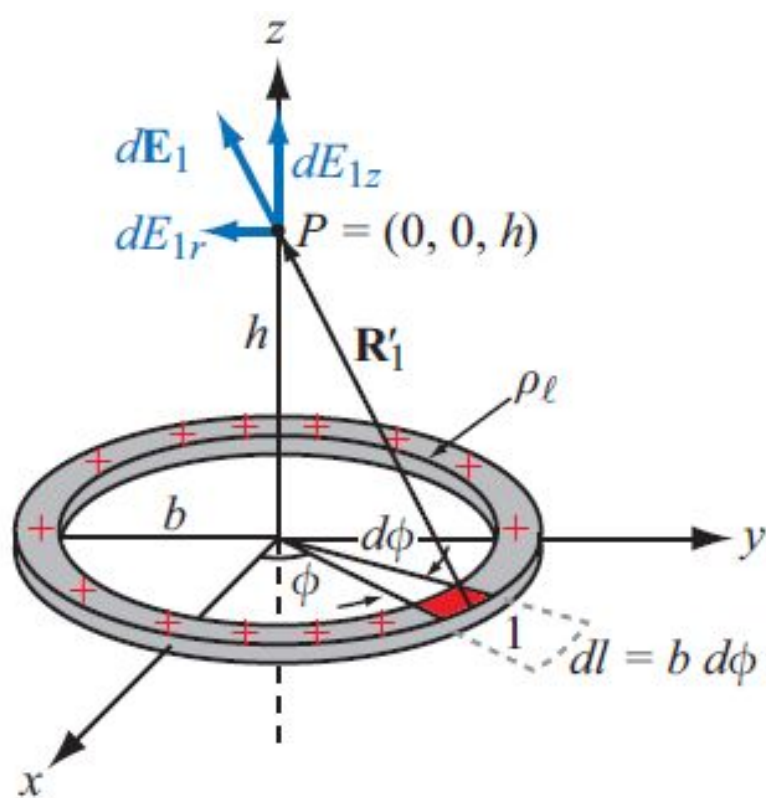
charge in this segment: $dq = \rho_{\ell} dl = \rho_{\ell} b d\phi$

vector from this segment to $(0, 0, h)$:

$$\mathbf{R}'_1 = -\hat{\mathbf{r}}b + \hat{\mathbf{z}}h$$

Example 4-4

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{l'} \hat{\mathbf{R}}' \frac{\rho_{\ell} dl'}{R'^2}$$



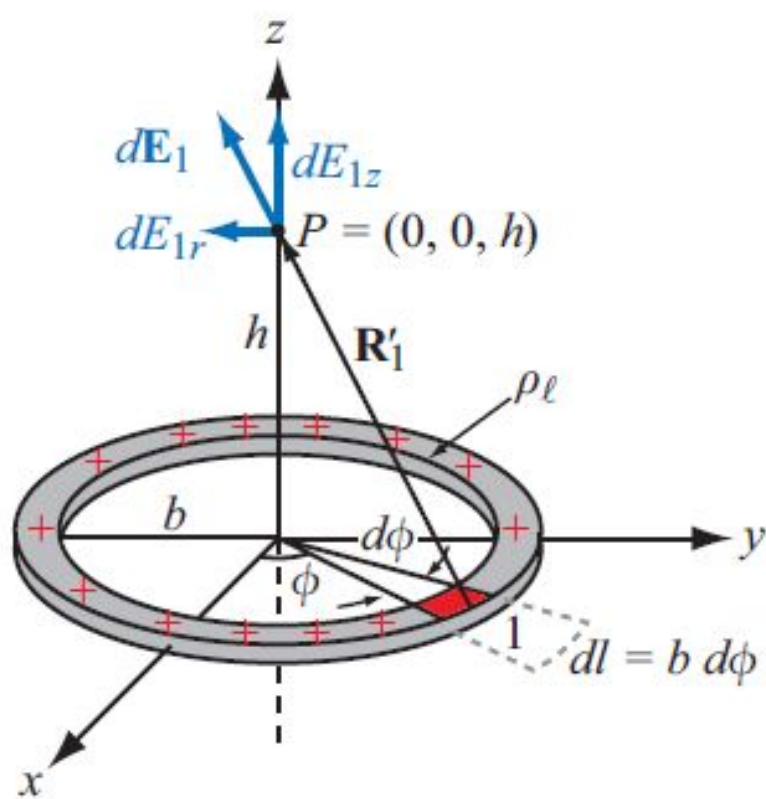
$$\mathbf{R}'_1 = -\hat{\mathbf{r}}b + \hat{\mathbf{z}}h,$$

from which it follows that

$$R'_1 = |\mathbf{R}'_1| = \sqrt{b^2 + h^2}, \quad \hat{\mathbf{R}}'_1 = \frac{\mathbf{R}'_1}{|\mathbf{R}'_1|} = \frac{-\hat{\mathbf{r}}b + \hat{\mathbf{z}}h}{\sqrt{b^2 + h^2}}.$$

Example 4-4

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{l'} \hat{\mathbf{R}}' \frac{\rho_{\ell} dl'}{R'^2}$$

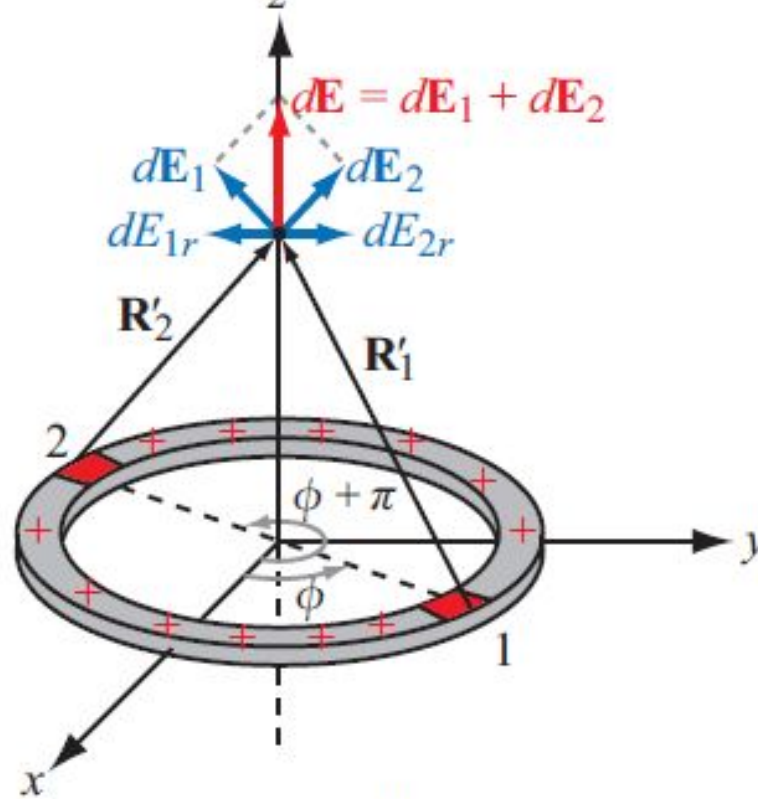


The electric field at $P = (0, 0, h)$ due to the charge in segment 1 therefore is

$$d\mathbf{E}_1 = \frac{1}{4\pi\epsilon_0} \hat{\mathbf{R}}'_1 \frac{\rho_{\ell} dl}{R'_1{}^2} = \frac{\rho_{\ell} b}{4\pi\epsilon_0} \frac{(-\hat{\mathbf{r}}b + \hat{\mathbf{z}}h)}{(b^2 + h^2)^{3/2}} d\phi.$$

Example 4-4

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{l'} \hat{\mathbf{R}}' \frac{\rho_l dl'}{R'^2}$$



Segment 1 has an \mathbf{E} field like: $-\hat{\mathbf{r}}b + \hat{\mathbf{z}}h$

and from symmetry, for segment 2: $+\hat{\mathbf{r}}b + \hat{\mathbf{z}}h$

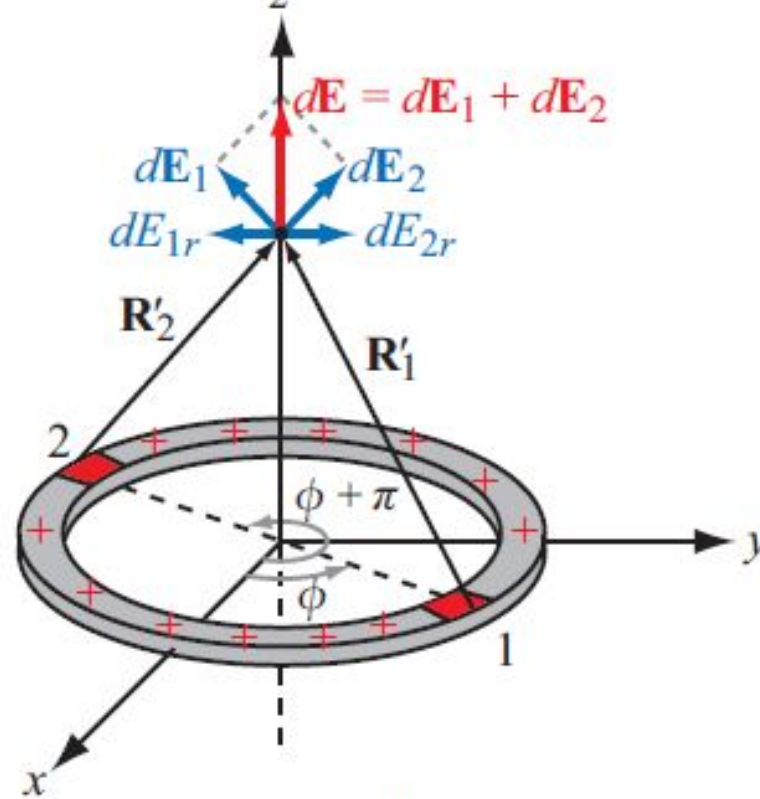
Therefore, their sum is:

$$d\mathbf{E} = d\mathbf{E}_1 + d\mathbf{E}_2 = \hat{\mathbf{z}} \frac{\rho_l b h}{2\pi\epsilon_0} \frac{d\phi}{(b^2 + h^2)^{3/2}}$$

Cont.

Example 4-4

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{l'} \hat{\mathbf{R}}' \frac{\rho_l dl'}{R'^2}$$



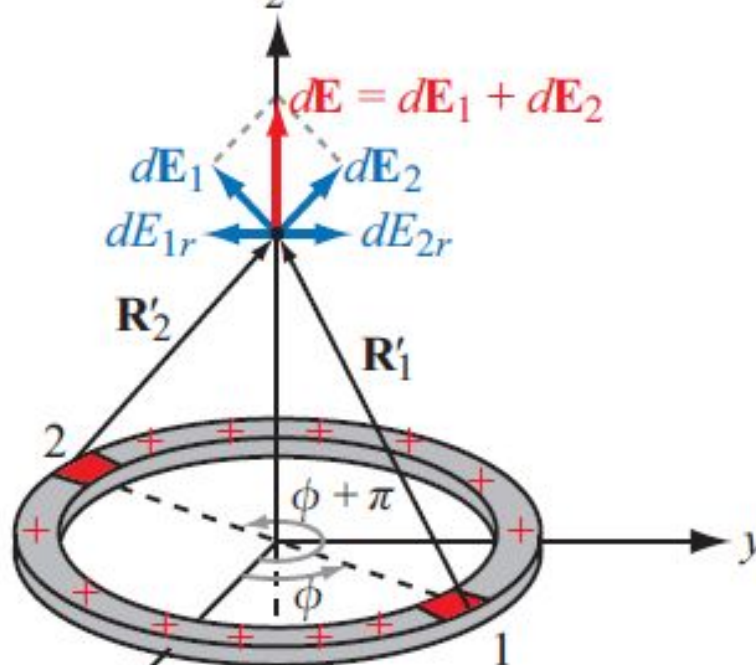
Since $d\mathbf{E}$ includes contributions from both halves of the ring,

We integrate $d\mathbf{E}$ around HALF the ring:

$$\mathbf{E} = \hat{\mathbf{z}} \frac{\rho_l b h}{2\pi\epsilon_0 (b^2 + h^2)^{3/2}} \int_0^\pi d\phi$$

Example 4-4

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{l'} \hat{\mathbf{R}}' \frac{\rho_{\ell} dl'}{R'^2}$$



$$\begin{aligned} \mathbf{E} &= \hat{\mathbf{z}} \frac{\rho_{\ell} b h}{2\pi\epsilon_0 (b^2 + h^2)^{3/2}} \int_0^{\pi} d\phi \\ &= \hat{\mathbf{z}} \frac{\rho_{\ell} b h}{2\epsilon_0 (b^2 + h^2)^{3/2}} \\ &= \hat{\mathbf{z}} \frac{h}{4\pi\epsilon_0 (b^2 + h^2)^{3/2}} Q, \end{aligned} \quad (4.23)$$

where $Q = 2\pi b\rho_{\ell}$ is the total charge on the ring.

Cont.

Example 4-4

Special cases:

$h=0$ (ring center):

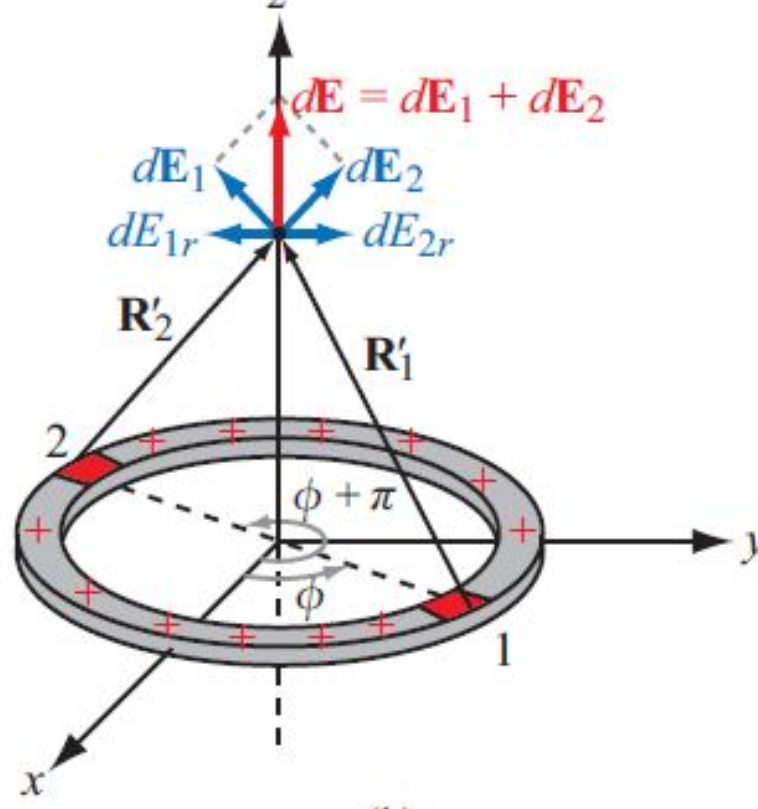
$$\mathbf{E} = 0$$

the effect of the charges cancels out

$h=\infty$:

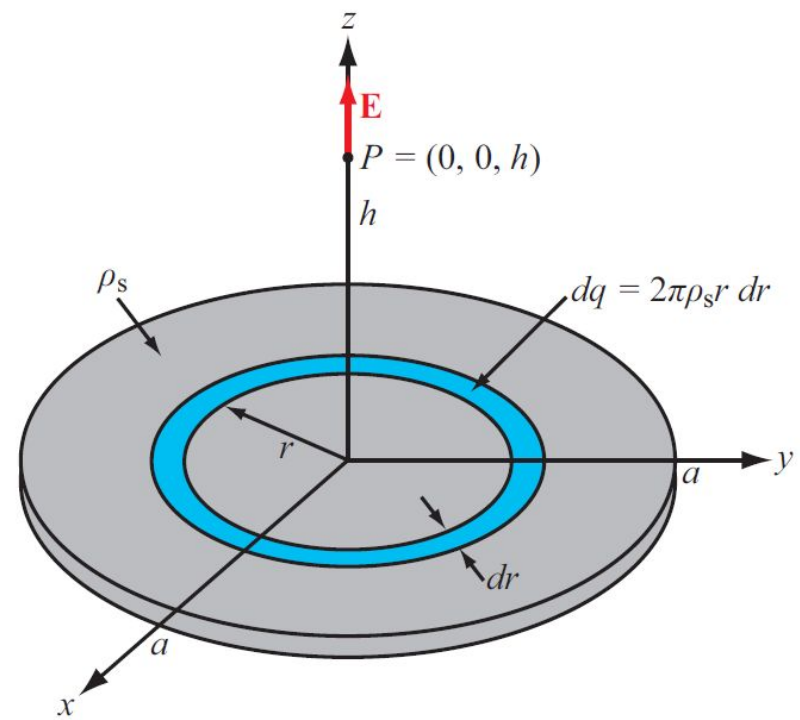
$$\mathbf{E} = \hat{\mathbf{z}} \frac{1}{4\pi\epsilon_0 h^2} Q$$

like a point charge at the origin



Example 4-5

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{S'} \hat{\mathbf{R}}' \frac{\rho_s ds'}{R'^2}$$

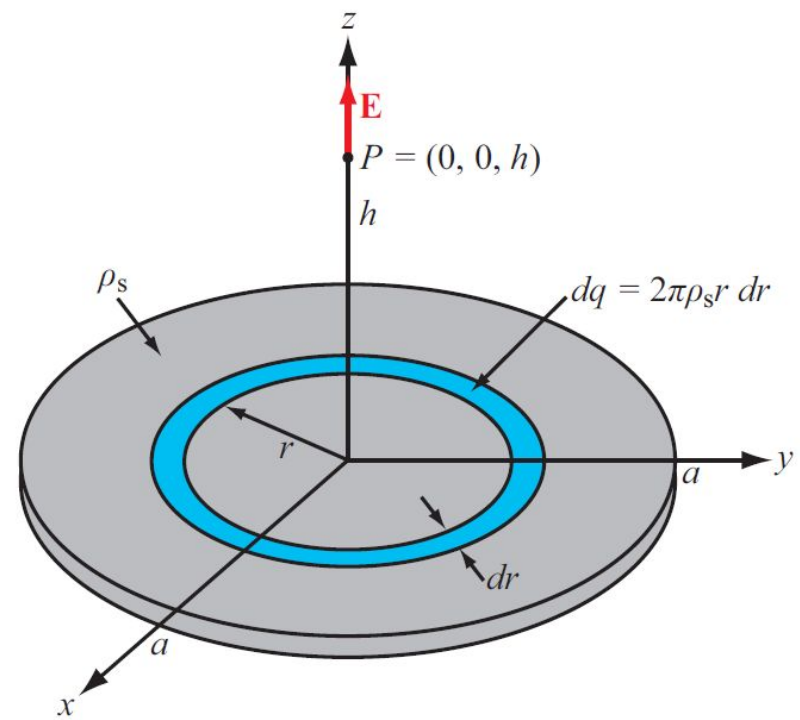


Given: Circular disk of radius: a
uniform charge density: ρ_s
in x-y plane

Find: $\mathbf{E}(0,0,h)$

Example 4-5

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{S'} \hat{\mathbf{R}}' \frac{\rho_s ds'}{R'^2}$$



Solution:

Use the solution to Example 4-4.

The sheet is a summation of concentric rings.
each ring of radius r , width dr :

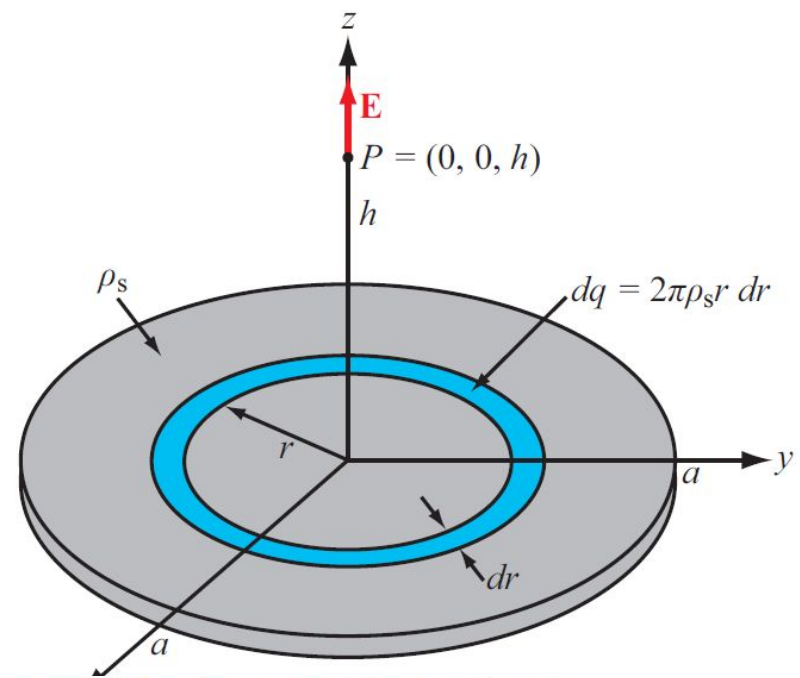
$$\text{area: } ds = 2\pi r dr$$

$$\text{charge: } dq = \rho_s ds = \rho_s 2\pi r dr$$

$$d\mathbf{E} = \hat{\mathbf{z}} \frac{h}{4\pi\epsilon_0 (r^2 + h^2)^{3/2}} (2\pi\rho_s r dr)$$

Example 4-5

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{S'} \hat{\mathbf{R}}' \frac{\rho_s ds'}{R'^2}$$

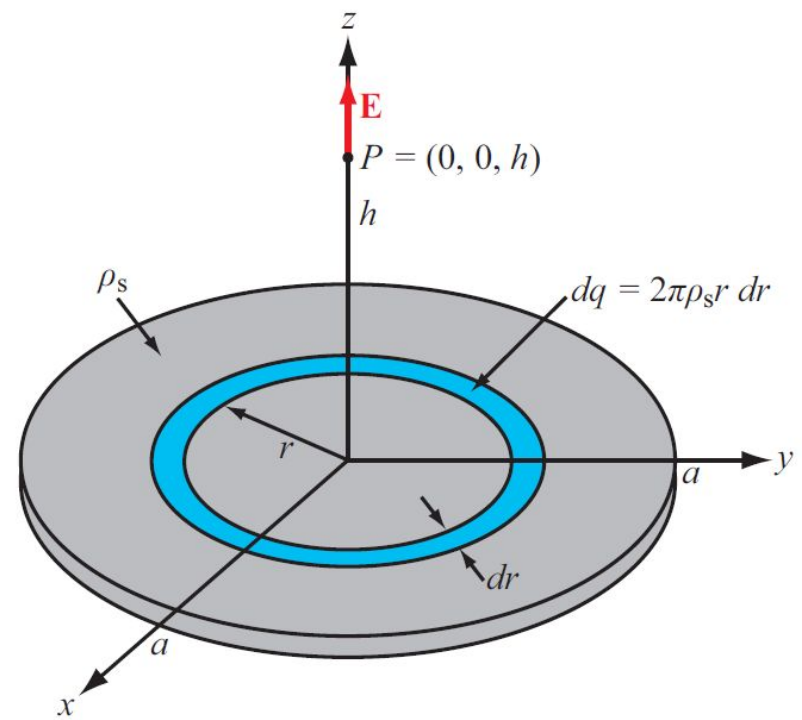


The total field at P is obtained by integrating the expression over the limits $r = 0$ to $r = a$:

$$\mathbf{E} = \hat{\mathbf{z}} \frac{\rho_s h}{2\epsilon_0} \int_0^a \frac{r dr}{(r^2 + h^2)^{3/2}}$$

Example 4-5

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{S'} \hat{\mathbf{R}}' \frac{\rho_s ds'}{R'^2}$$

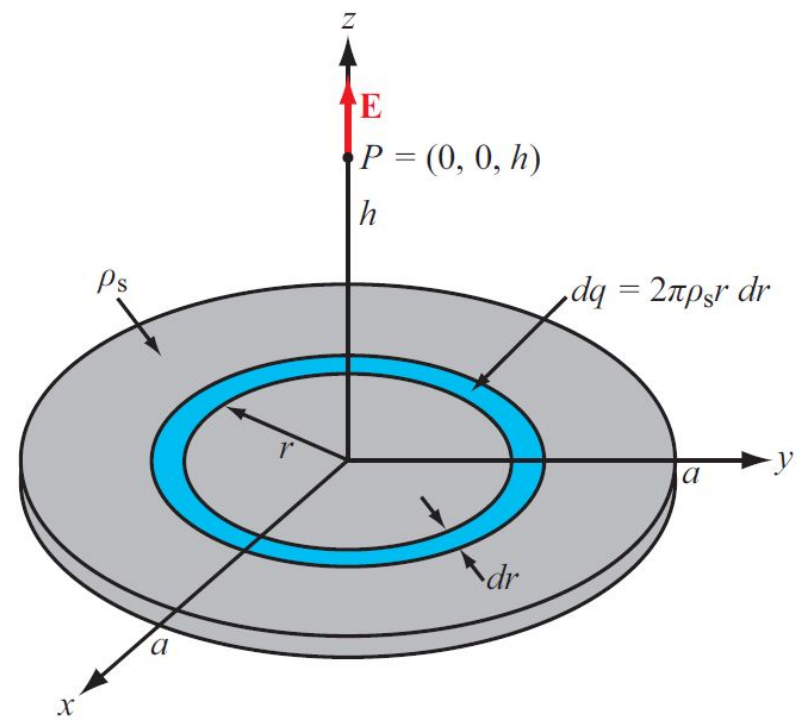


$$\int \frac{x dx}{(x^2 + c^2)^{3/2}} = -\frac{1}{(x^2 + c^2)^{1/2}}$$

$$\int_0^a \frac{r dr}{(r^2 + h^2)^{3/2}} = -\frac{1}{(r^2 + h^2)^{1/2}} \Big|_0^a$$

Example 4-5

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{S'} \hat{\mathbf{R}}' \frac{\rho_s ds'}{R'^2}$$



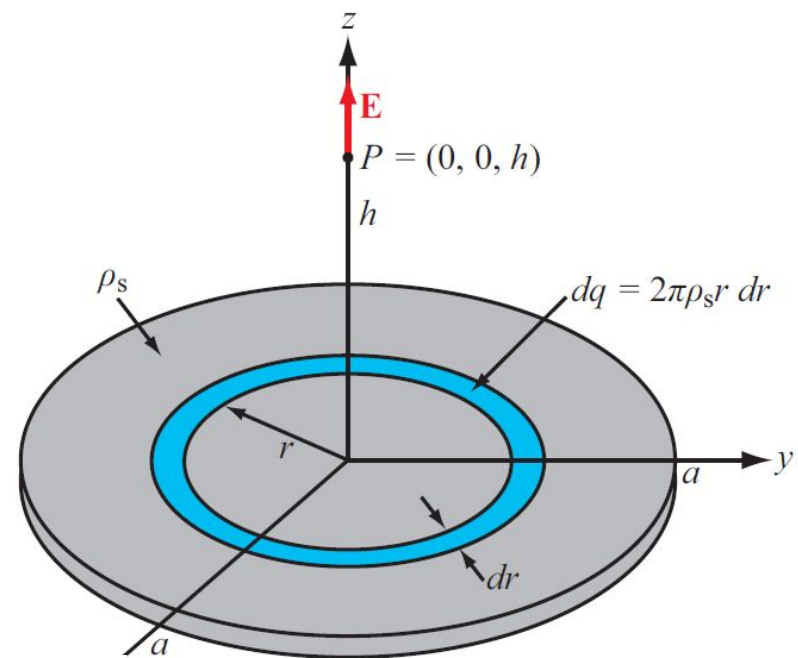
Using the + square-root:

$$-\frac{1}{(r^2 + h^2)^{1/2}} \Big|_{r=0}^a = -\frac{1}{(a^2 + h^2)^{1/2}} + \frac{1}{(h^2)^{1/2}}$$

$$-\frac{1}{(r^2 + h^2)^{1/2}} \Big|_{r=0}^a = -\frac{1}{(a^2 + h^2)^{1/2}} + \frac{1}{|h|}$$

Example 4-5

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{S'} \hat{\mathbf{R}}' \frac{\rho_s ds'}{R'^2}$$



$$\mathbf{E} = \hat{\mathbf{z}} \frac{\rho_s h}{2\epsilon_0} \left[\frac{1}{|h|} - \frac{1}{(a^2 + h^2)^{1/2}} \right]$$

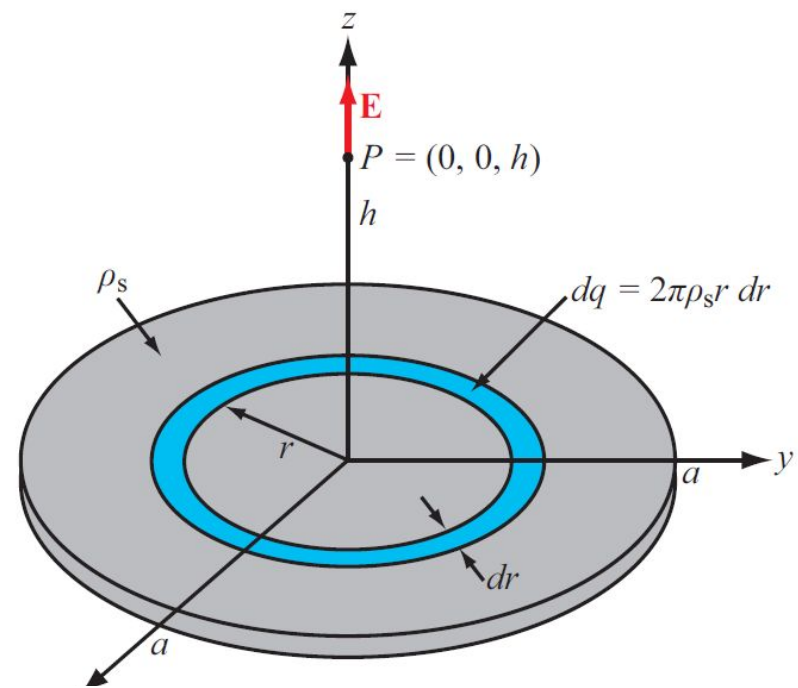
$$\mathbf{E} = \hat{\mathbf{z}} \frac{\rho_s}{2\epsilon_0} \left[\frac{h}{|h|} - \frac{h}{(a^2 + h^2)^{1/2}} \right]$$

$$\mathbf{E} = \hat{\mathbf{z}} \frac{\rho_s}{2\epsilon_0} \left[1 - \frac{h}{(a^2 + h^2)^{1/2}} \right] \quad (h > 0)$$

$$\mathbf{E} = \hat{\mathbf{z}} \frac{\rho_s}{2\epsilon_0} \left[-1 - \frac{h}{(a^2 + h^2)^{1/2}} \right] \quad (h < 0)$$

Example 4-5

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{S'} \hat{\mathbf{R}}' \frac{\rho_s ds'}{R'^2}$$



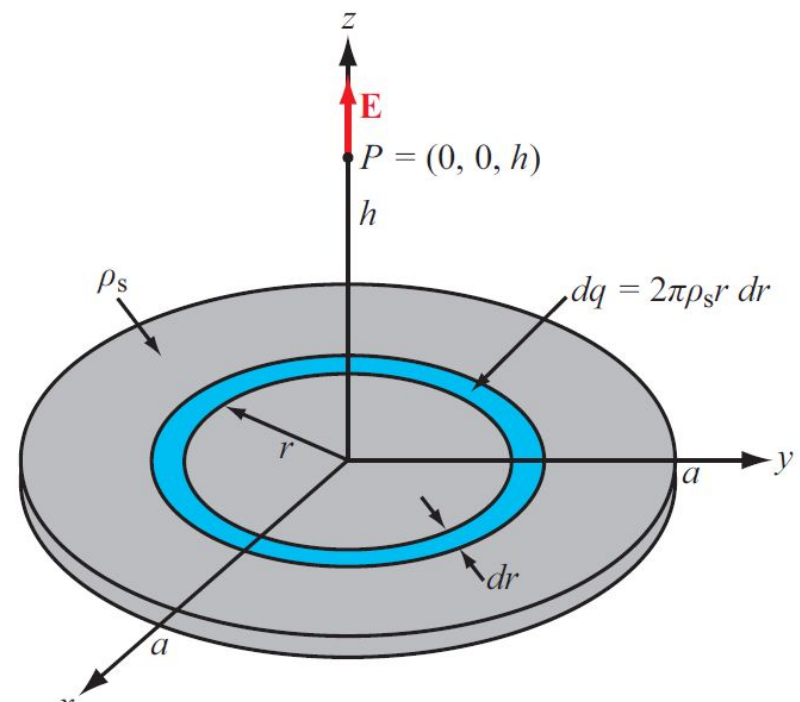
$$\mathbf{E} = \hat{\mathbf{z}} \frac{\rho_s}{2\epsilon_0} \left[-1 - \frac{h}{(a^2 + h^2)^{1/2}} \right] \quad (h < 0)$$

$$\mathbf{E} = \hat{\mathbf{z}} \frac{\rho_s}{2\epsilon_0} \left[-1 + \frac{|h|}{(a^2 + h^2)^{1/2}} \right] \quad (h < 0)$$

$$\mathbf{E} = -\hat{\mathbf{z}} \frac{\rho_s}{2\epsilon_0} \left[1 - \frac{|h|}{(a^2 + h^2)^{1/2}} \right] \quad (h < 0)$$

Example 4-5

$$\mathbf{E} = \frac{1}{4\pi\epsilon} \int_{S'} \hat{\mathbf{R}}' \frac{\rho_s ds'}{R'^2}$$



$$\mathbf{E} = -\hat{\mathbf{z}} \frac{\rho_s}{2\epsilon_0} \left[1 - \frac{|h|}{(a^2 + h^2)^{1/2}} \right] \quad (h < 0)$$

$$\mathbf{E} = +\hat{\mathbf{z}} \frac{\rho_s}{2\epsilon_0} \left[1 - \frac{|h|}{(a^2 + h^2)^{1/2}} \right] \quad (h > 0)$$

$$\mathbf{E} = \pm \hat{\mathbf{z}} \frac{\rho_s}{2\epsilon_0} \left[1 - \frac{|h|}{(a^2 + h^2)^{1/2}} \right] \quad (+ \text{ if } h > 0, - \text{ if } h < 0)$$

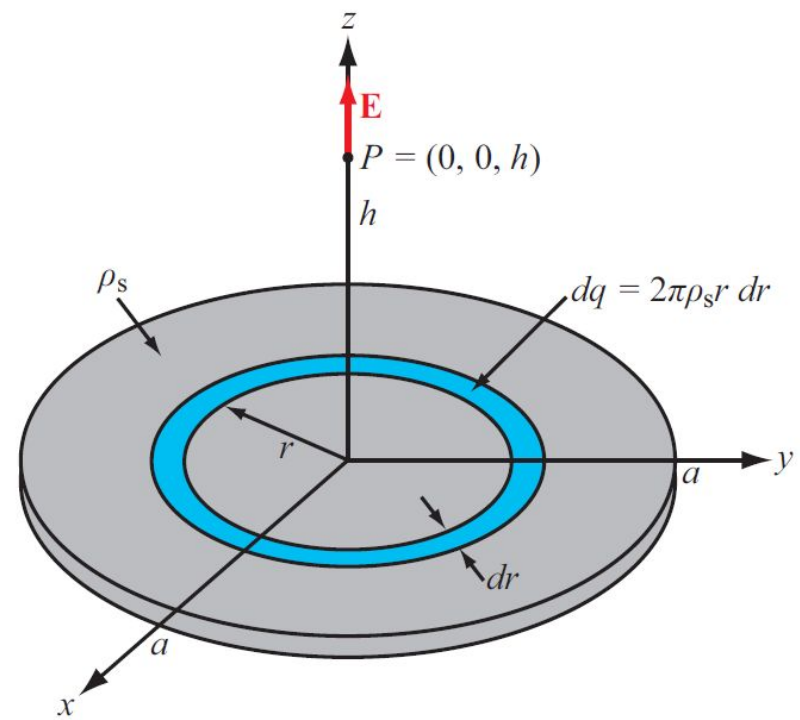
Example 4-5

Special case:

$$\mathbf{E} := \pm \hat{\mathbf{z}} \frac{\rho_s}{2\epsilon_0} \left[1 - \frac{|h|}{\sqrt{a^2 + h^2}} \right]$$

$a = \infty$ (infinite sheet)

$$\mathbf{E} = \pm \hat{\mathbf{z}} \frac{\rho_s}{2\epsilon_0} \quad (\text{infinite sheet of charge}).$$

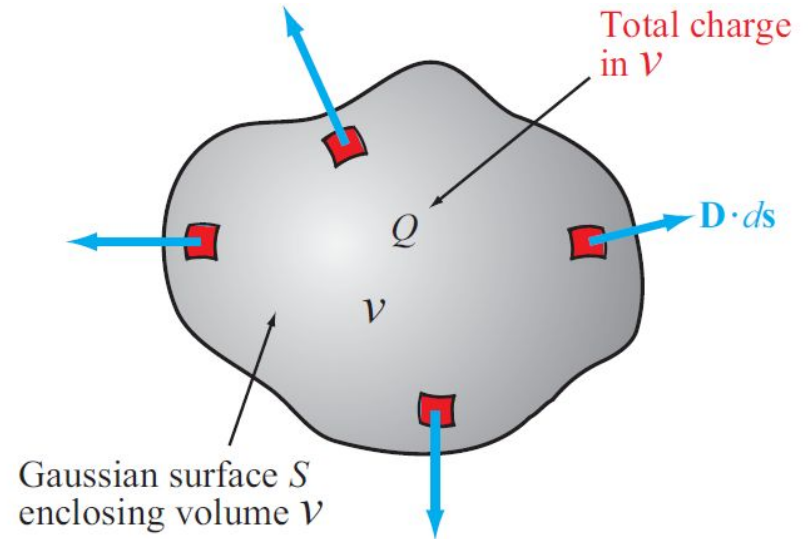


constant: does not depend on distance!

4-4 Gauss's Law

$$\nabla \cdot \mathbf{D} = \rho_v$$

(Differential form of Gauss's law),



4-4 Gauss's Law

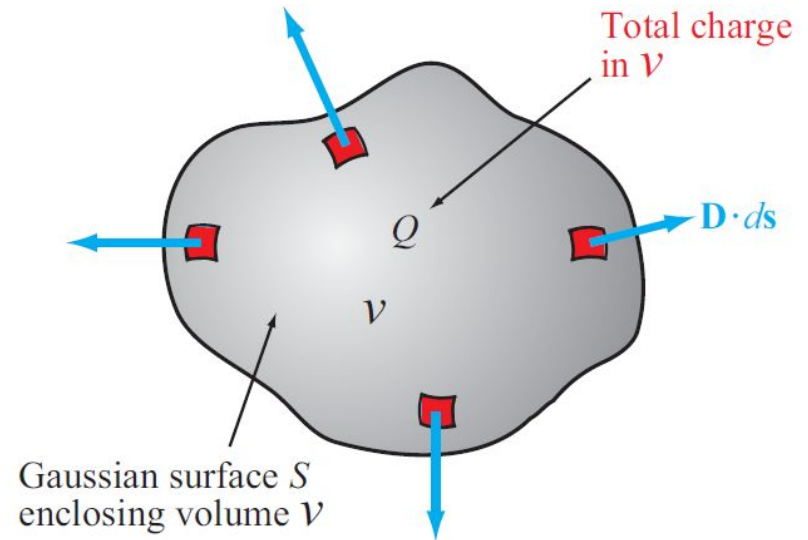
Application of the divergence theorem gives:

$$\int_V \nabla \cdot \mathbf{D} dV = \oint_S \mathbf{D} \cdot d\mathbf{s}.$$

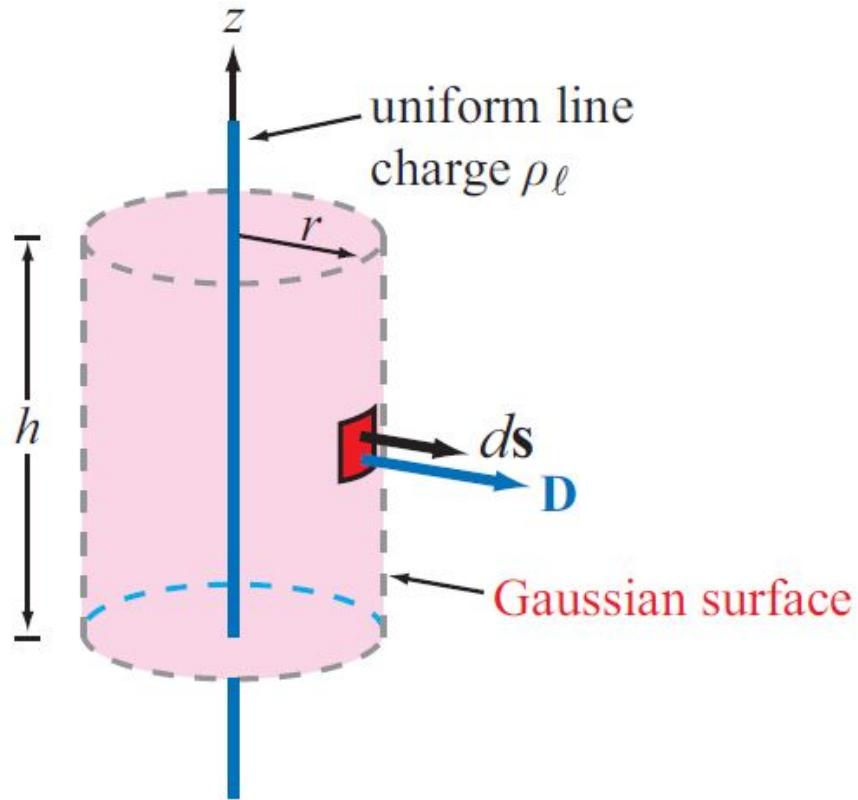
$$\int_V \rho_v dV = \oint_S \mathbf{D} \cdot d\mathbf{s}$$

$$\oint_S \mathbf{D} \cdot d\mathbf{s} = Q$$

(Integral form of Gauss's law).



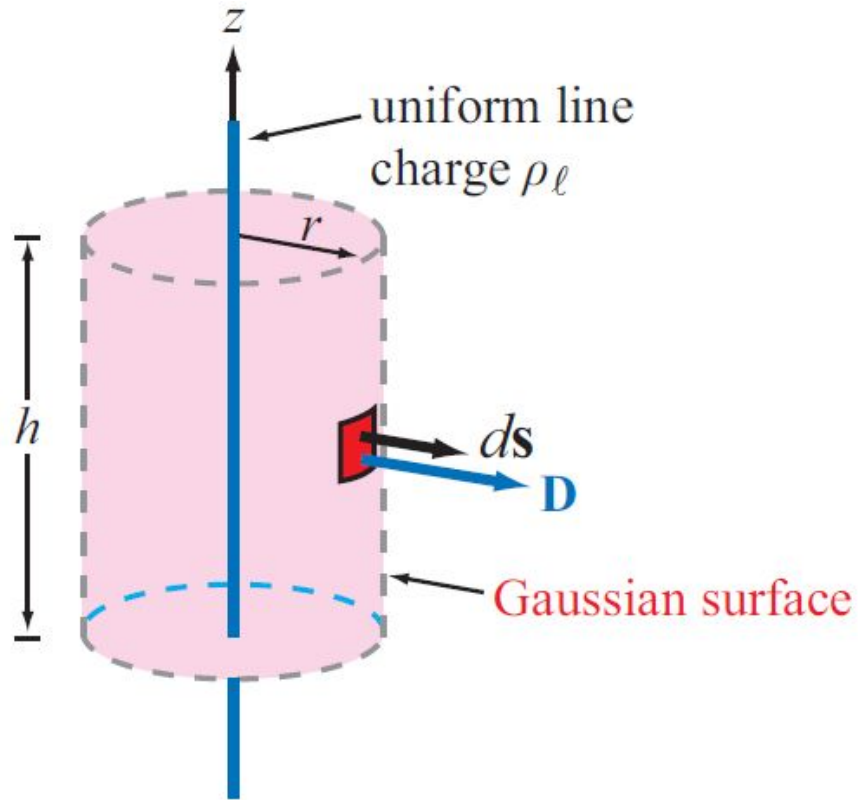
Example 4-6



Given: line charge density: ρ_ℓ
infinite
along z -axis
in free space

Find: Using Gauss's Law:
 \mathbf{E} everywhere

Example 4-6



(Don't make it infinite, since that often causes problems...)

Solution:

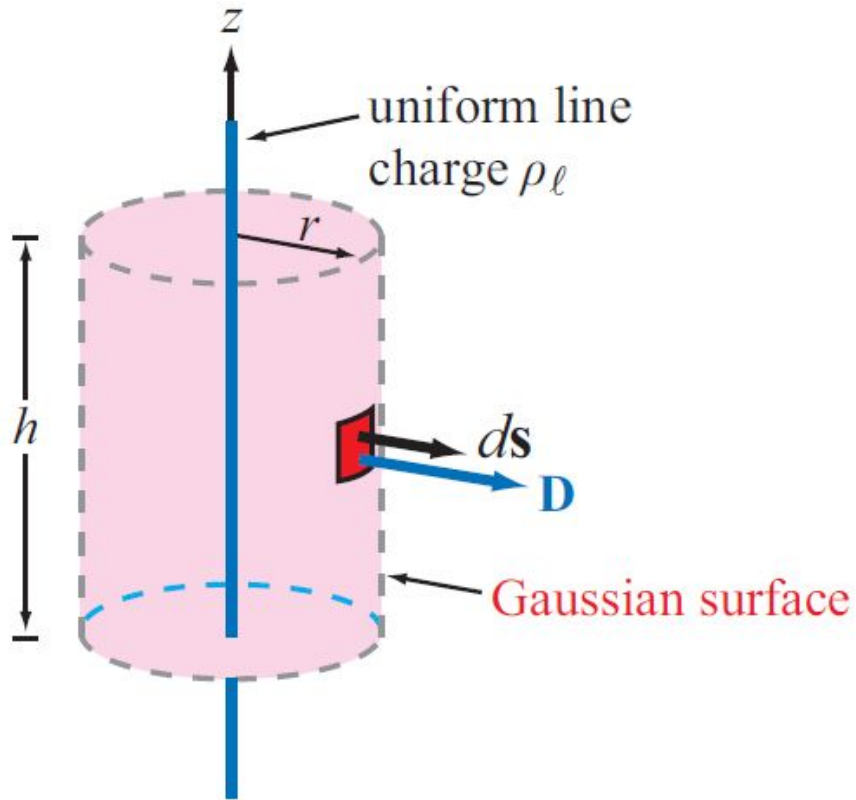
Since we are using Gauss's law:

$$\oint_S \mathbf{D} \cdot d\mathbf{s} = Q.$$

we need to identify an appropriate surface.

Construct an imaginary cylinder of radius r and height h

Example 4-6



Note that we will integrate over the curved surface of the cylinder, and not include the 2 flat ends.

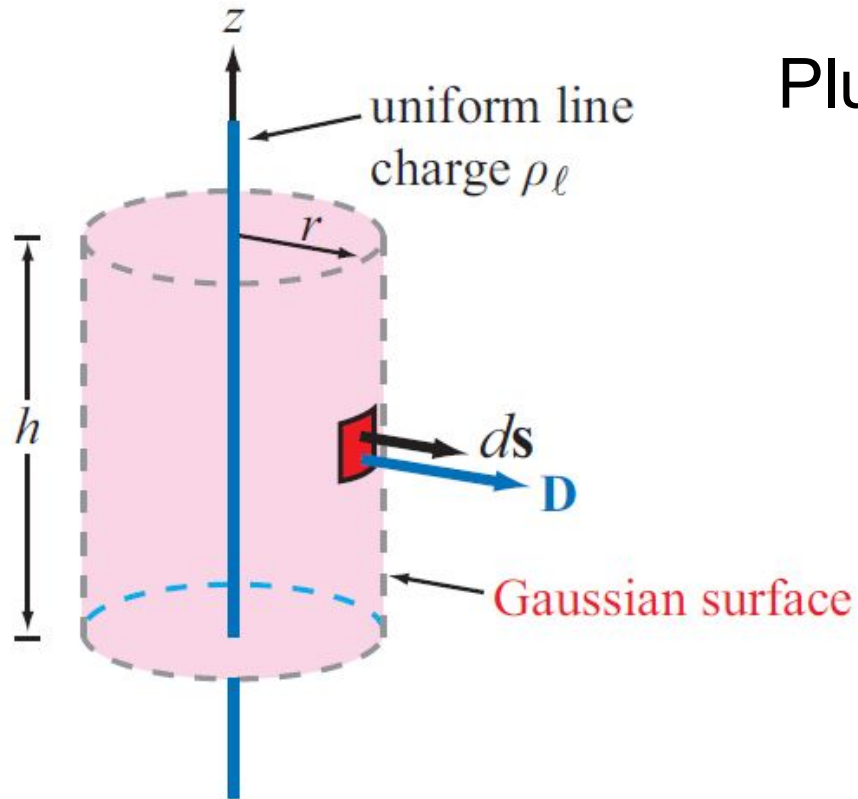
ds is surface patch on this curved surface:

direction: $\hat{\mathbf{r}}$

magnitude: $r d\phi dz$

\mathbf{D} is radial: $\hat{\mathbf{r}} D_r$

Example 4-6



Plug in:

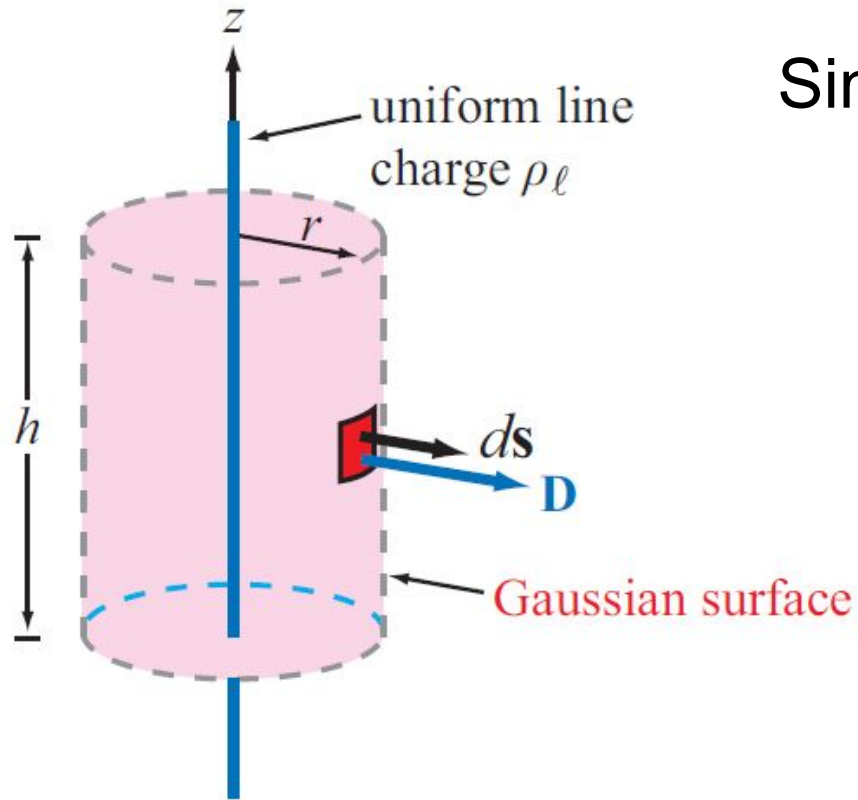
$$\int_{z=0}^h \int_{\phi=0}^{2\pi} \mathbf{D} \cdot \hat{\mathbf{r}} r d\phi dz = \rho_l h$$

$$\int_{z=0}^h \int_{\phi=0}^{2\pi} \hat{\mathbf{r}} D_r \cdot \hat{\mathbf{r}} r d\phi dz = \rho_l h$$

$$D_r r \int_{z=0}^h \int_{\phi=0}^{2\pi} d\phi dz = \rho_l h$$

$$D_r r \int_{z=0}^h dz \int_{\phi=0}^{2\pi} d\phi = \rho_l h$$

Example 4-6



Simplify:

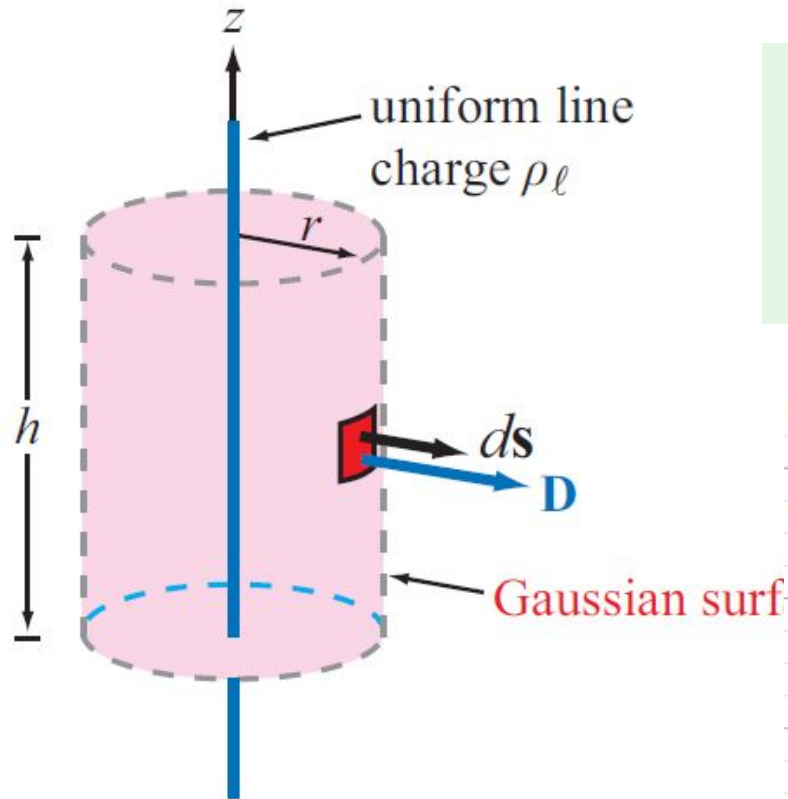
$$D_r r \int_{z=0}^h \int_{\phi=0}^{2\pi} d\phi dz = \rho_l h$$

$$D_r r \int_{z=0}^h dz \int_{\phi=0}^{2\pi} d\phi = \rho_l h$$

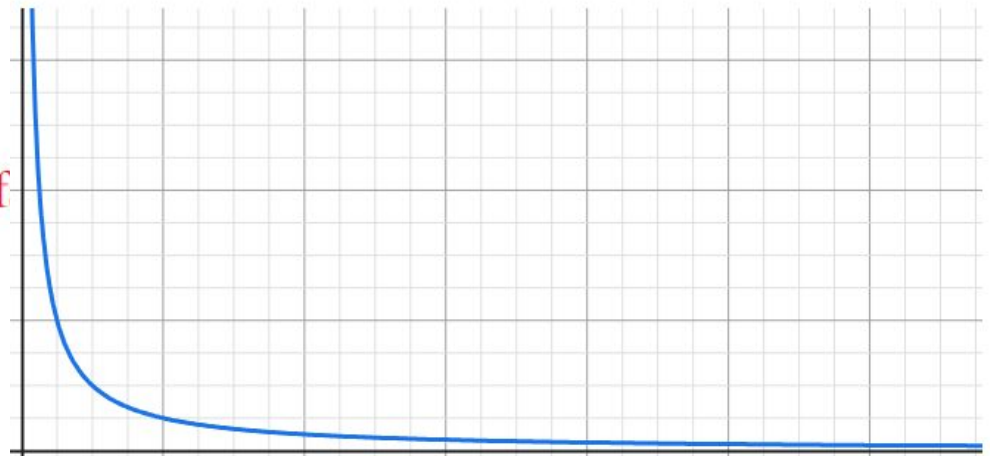
$$D_r r \left[z \right]_{z=0}^h \left[\phi \right]_{\phi=0}^{2\pi} = \rho_l h$$

$$D_r r h 2\pi = \rho_l h$$

Example 4-6

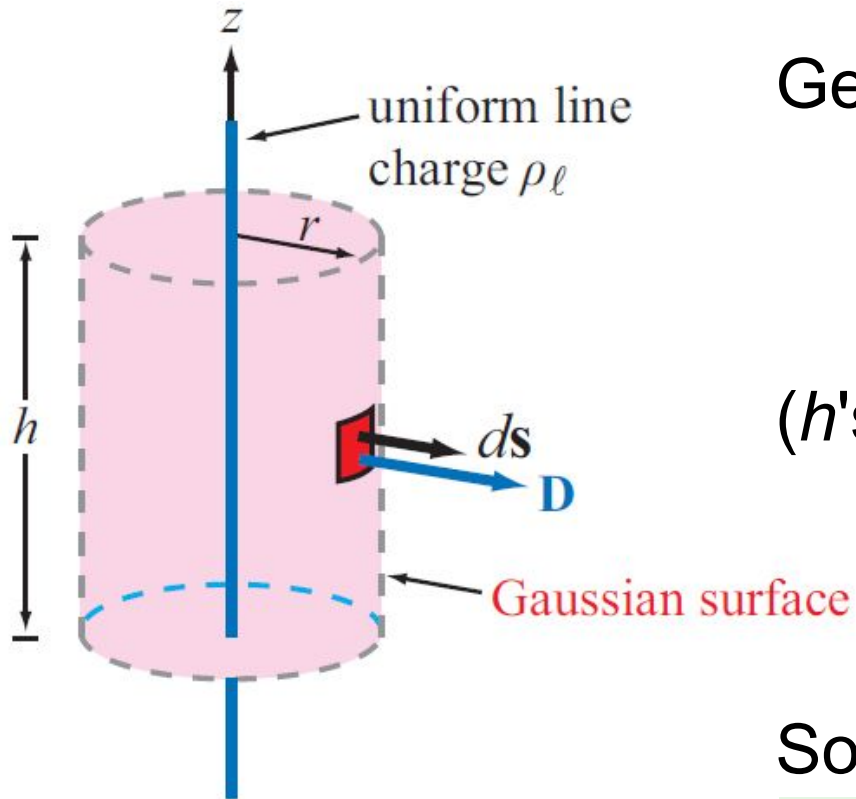


$$\mathbf{E} = \frac{\mathbf{D}}{\epsilon_0} = \hat{\mathbf{r}} \frac{D_r}{\epsilon_0} = \hat{\mathbf{r}} \frac{\rho_\ell}{2\pi\epsilon_0 r}$$



Note: invalid for r too small.

Example 4-6



Get:

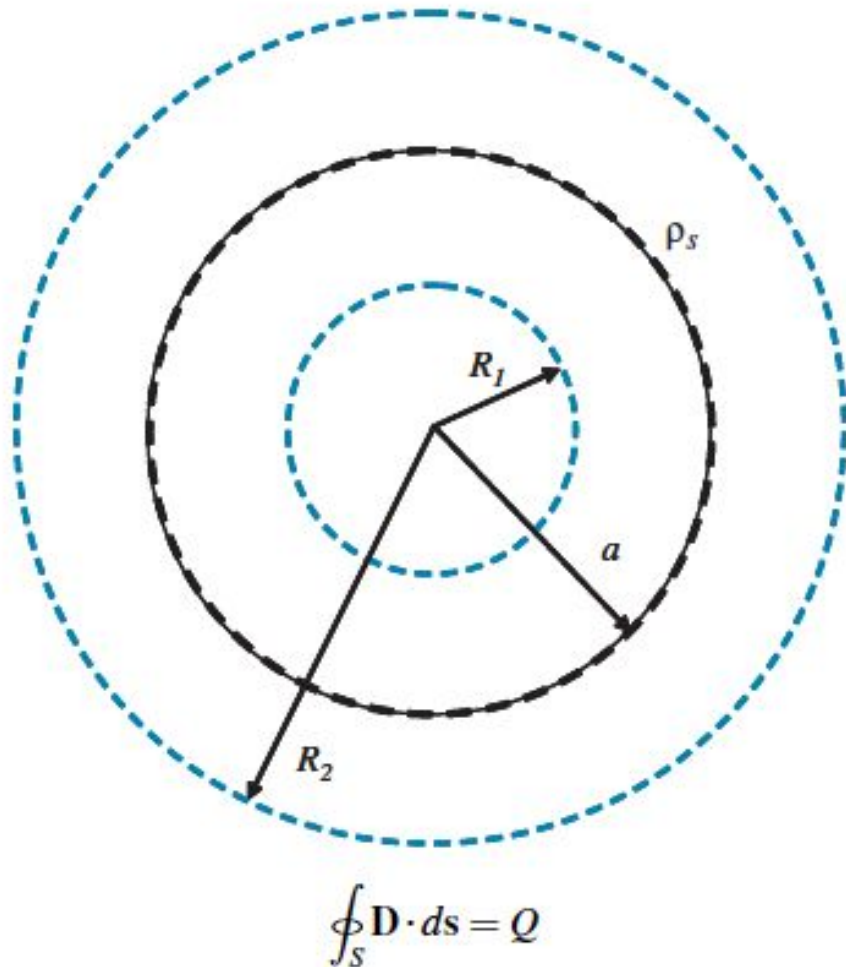
$$D_r = \frac{\rho_\ell}{2\pi r}$$

(h 's cancel)

Solving for \mathbf{E} :

$$\mathbf{E} = \frac{\mathbf{D}}{\epsilon_0} = \hat{\mathbf{r}} \frac{D_r}{\epsilon_0} = \hat{\mathbf{r}} \frac{\rho_\ell}{2\pi\epsilon_0 r}$$

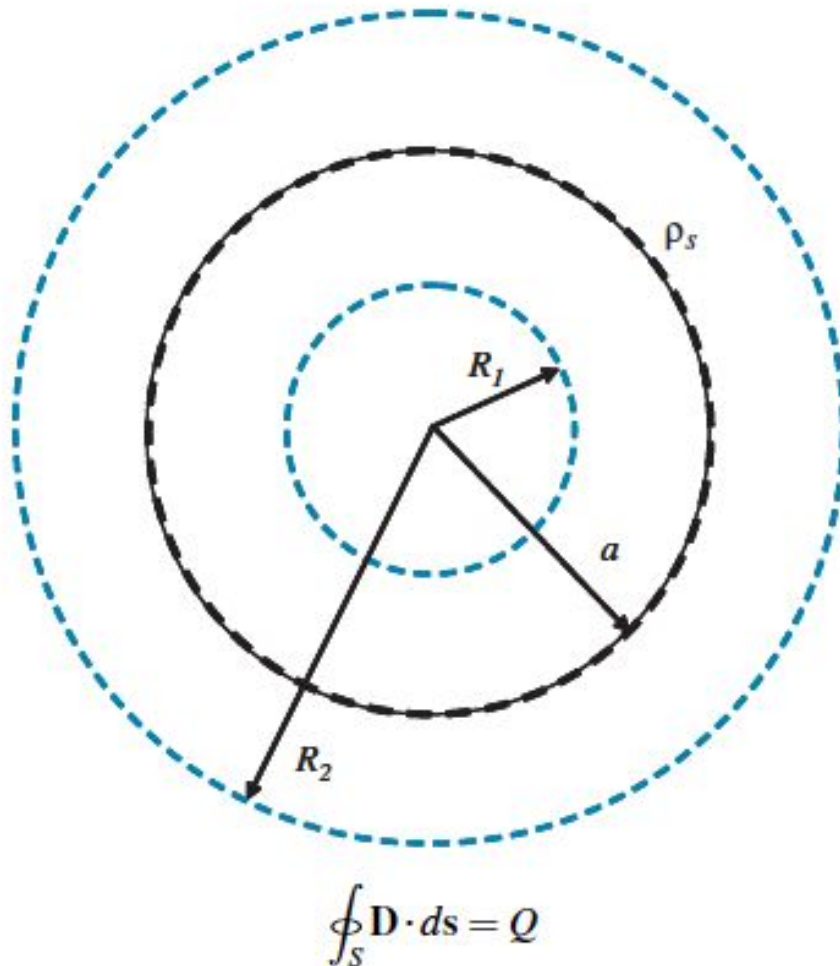
Exercise 4-8



Given: surface charge density: ρ_s
on sphere, radius a
in some medium with ϵ

Find: Using Gauss's Law:
 \mathbf{E} everywhere
(both inside and outside the
spherical shell)

Exercise 4-8



Solution:

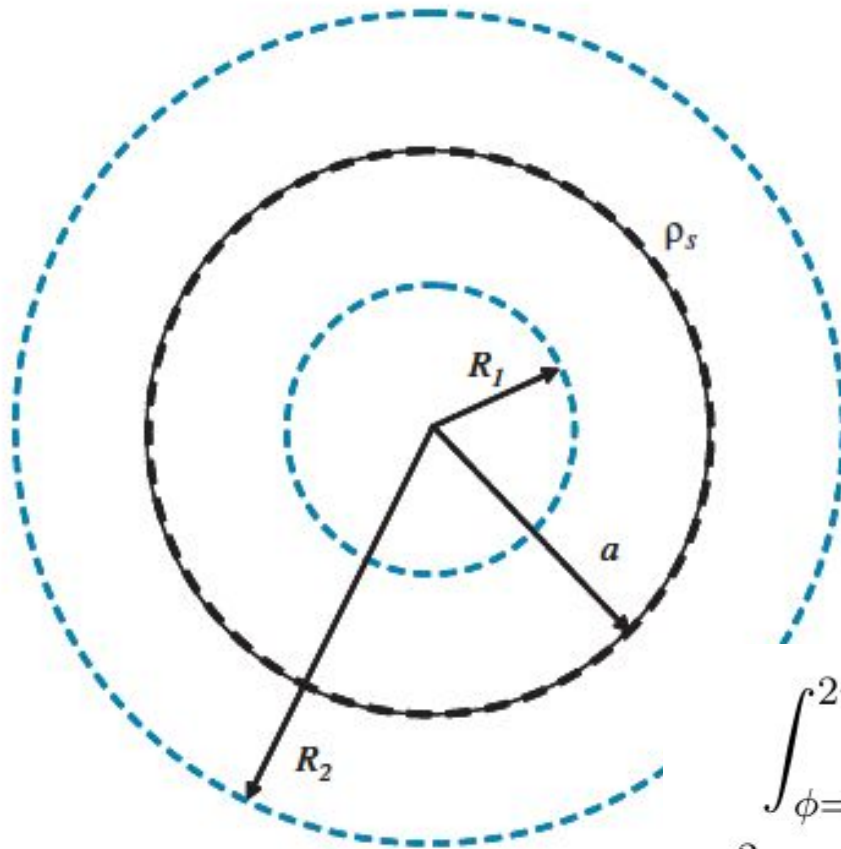
Since we are using Gauss's law:

$$\oint_S \mathbf{D} \cdot d\mathbf{s} = Q.$$

we need to identify an appropriate surface.

Choose a spherical surface of radius $R_1 < a$ for fields inside the sphere, and $R_2 > a$ for fields outside the sphere.

Exercise 4-8



$$\oint_S \mathbf{D} \cdot d\mathbf{s} = Q$$

$d\mathbf{s}$ is surface patch on this spherical surface:

direction: $\hat{\mathbf{R}}$

magnitude: $R^2 \sin \theta d\theta d\phi$

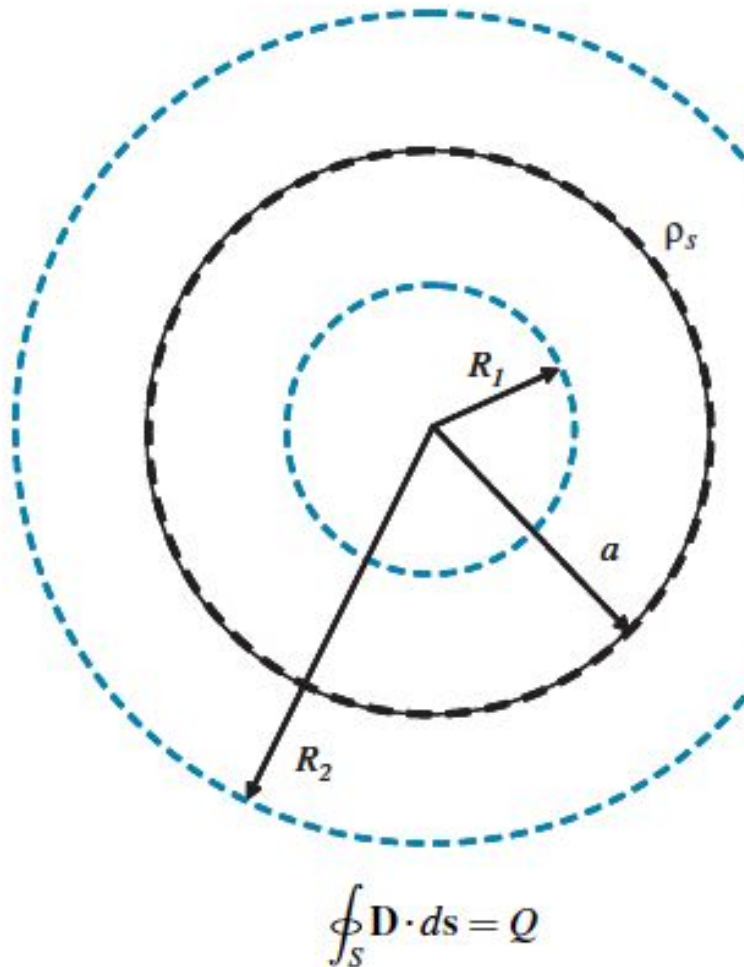
\mathbf{D} is radial

Plug in:

$$\int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \mathbf{D} \cdot \hat{\mathbf{R}} R^2 \sin \theta d\theta d\phi = Q_{\text{enclosed}}$$

$$\int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \hat{\mathbf{R}} D_R \cdot \hat{\mathbf{R}} R^2 \sin \theta d\theta d\phi = Q_{\text{enclosed}}$$

Exercise 4-8



Simplify:

$$\int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \hat{\mathbf{R}} D_R \cdot \hat{\mathbf{R}} R^2 \sin \theta d\theta d\phi = Q_{\text{enclosed}}$$

$$D_R R^2 \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \sin \theta d\theta d\phi = Q_{\text{enclosed}}$$

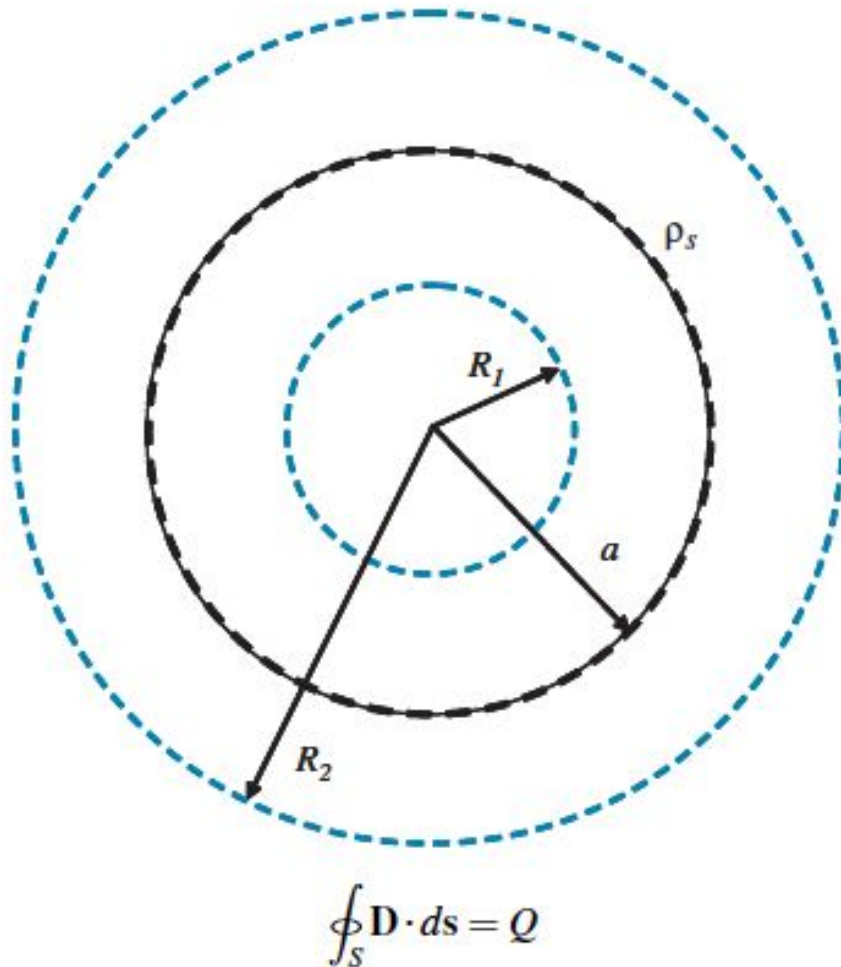
$$D_R R^2 \int_{\phi=0}^{2\pi} d\phi \int_{\theta=0}^{\pi} \sin \theta d\theta = Q_{\text{enclosed}}$$

$$D_R R^2 2\pi \left[-\cos \theta \right]_{\theta=0}^{\pi} = Q_{\text{enclosed}}$$

$$D_R R^2 2\pi 2 = Q_{\text{enclosed}}$$

$$D_R = \frac{Q_{\text{enclosed}}}{4\pi R^2}$$

Exercise 4-8



For the surface inside the spherical shell, $Q_{\text{enclosed}} = 0$, so $\mathbf{D} = 0$, and $\mathbf{E} = 0$.

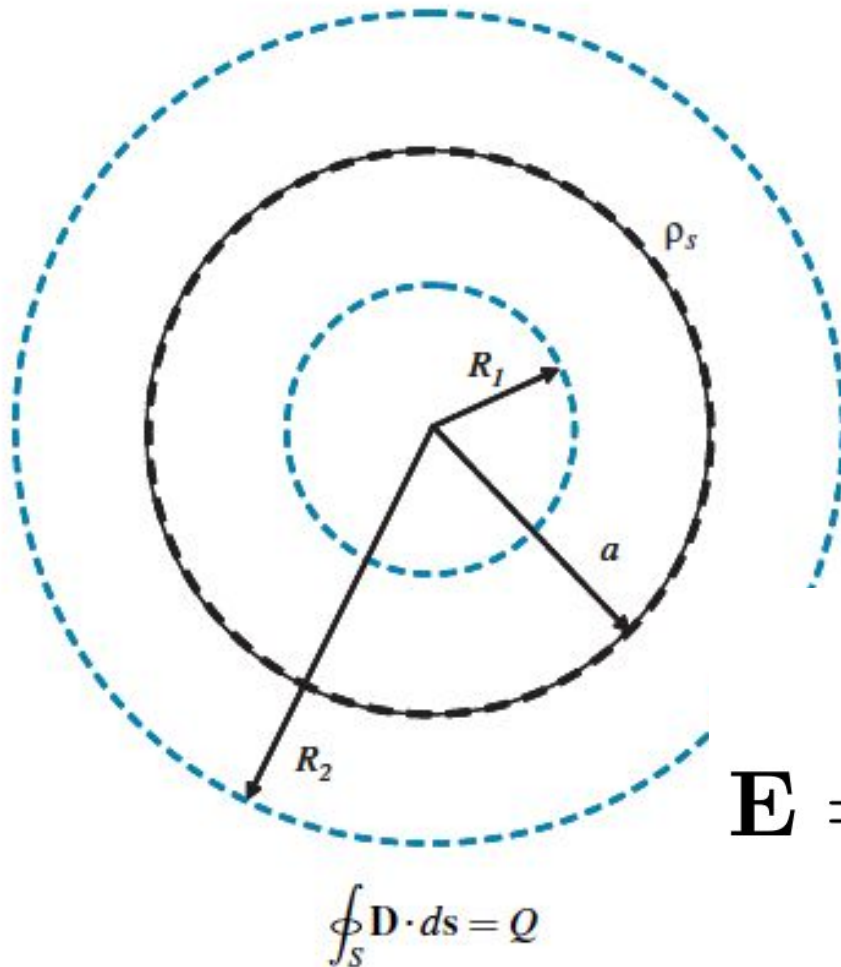
For the surface outside the spherical shell:

$$Q_{\text{enclosed}} = \rho_s A = \rho_s 4\pi a^2$$

So:

$$D_R = \frac{\rho_s 4\pi a^2}{4\pi R^2} = \frac{\rho_s a^2}{R^2}$$

Exercise 4-8



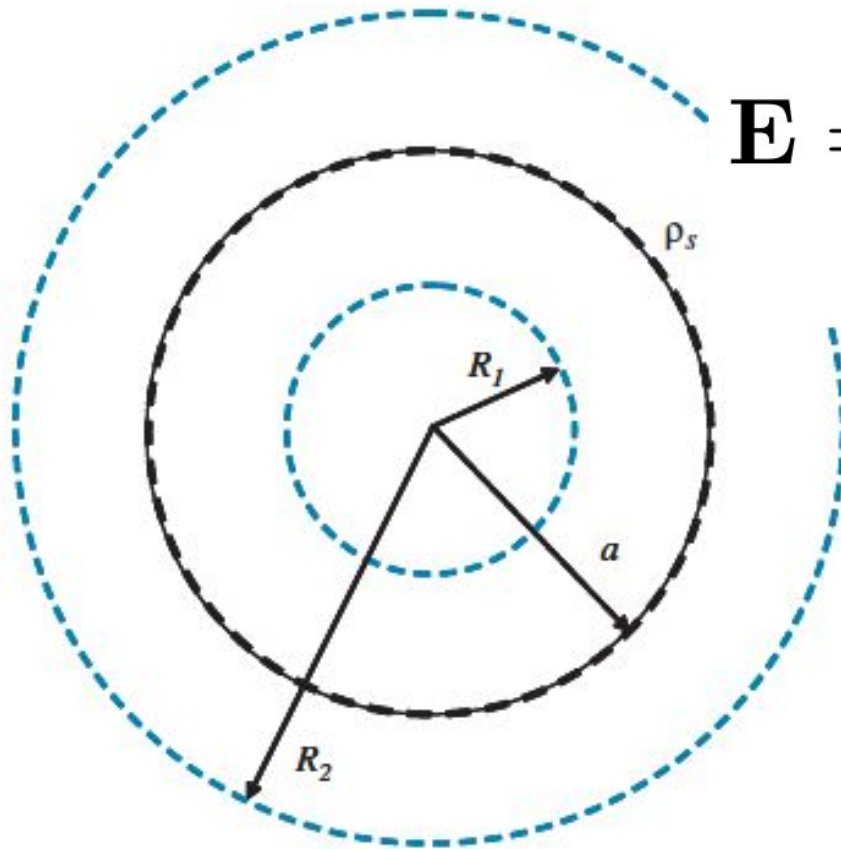
For outside the spherical shell:

$$\mathbf{E} = \frac{\mathbf{D}}{\epsilon} = \frac{\hat{\mathbf{R}}}{\epsilon} D_r = \hat{\mathbf{R}} \frac{\rho_s a^2}{\epsilon R^2}$$

The complete solution:

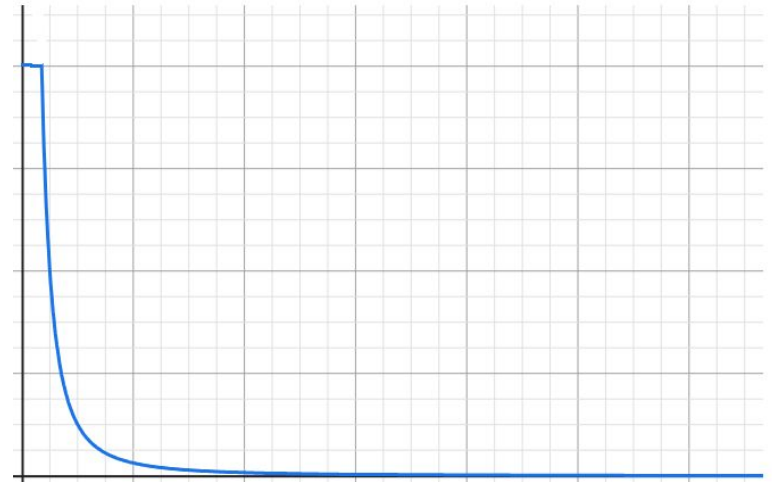
$$\mathbf{E} = \begin{cases} 0 & : R < a \\ \hat{\mathbf{R}} \frac{\rho_s a^2}{\epsilon R^2} & : R > a \end{cases}$$

Exercise 4-8



$$\mathbf{E} = \begin{cases} 0 & : R < a \\ \hat{\mathbf{R}} \frac{\rho_s a^2}{\epsilon R^2} & : R > a \end{cases}$$

$$\oint_S \mathbf{D} \cdot d\mathbf{s} = Q$$



Homework

76

Homework 12 is due tomorrow at midnight.

submit to gradescope via the canvas site.

Next Time



Section 4-5:

Electric Potential