## NONLIOCALITY IN SHALLOW QUANTUM CIRCUITS

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## INTRODUCTION

Quantum supremacy is expected to drastically change modern computing, but physical implementation is difficult due to nature of qubits.

Due to limitations, there is increasing interest in Noisy Intermediate-Scale Quantum technology (50-100 qubits).

In 2017, Bravyi, Gosset, and König publish "Quantum Advantage of Shallow Circuits".

Authors prove quantum circuits are more powerful than classical ones, without complexity theory assumptions. In particular,
$2 D-H L F P \in Q N C^{0} \backslash N C^{0}$.
We illustrated the importance of the nonlocality property in shallow quantum circuits by proving some of their results for some small examples.

$\mathrm{NC}^{\mathbf{0}}=$ complexity class of all poly-size circuits of $O(1)$ depth with bounded fan-in gates
$\mathrm{QNC}^{0}=$ quantum analog of $N C^{0}$

## Hidden Linear Function PROBLEM

Hidden Linear Function Problem (HLFP): Given: quadratic form $q: \mathbb{F}_{2}^{n} \rightarrow \mathbb{Z}_{4}$,

$$
q(x)=2 \sum_{1 \leq i<j \leq n} A_{i j} x_{i} x_{j}+\sum_{k=1}^{n} b_{k} x_{k},
$$ where $x_{k}, A_{i j}, b_{k} \in \mathbb{F}_{2}^{n}$.

Find: $z \in \mathbb{F}_{2}^{n}$ such that $q(x)=2 z^{T} x$, $\forall x \in \mathcal{L}_{q}$, where $\mathcal{L}_{q}$ denotes the set
$\left\{x \in \mathbb{F}_{2}^{n} \mid q(x \oplus y)=q(x)+q(y), \forall y \in \mathbb{F}_{2}^{n}\right\}$.

## 2D-HLFP:

Let $G=(V, E)$ denote the $N \times N$ grid graph and $A \in\{0,1\}^{|E|}$ be its adjacency matrix. Given: quadratic form $q: \mathbb{F}_{2}^{|V|} \rightarrow \mathbb{Z}_{4}$

$$
q(x)=2 \sum_{(u, v) \in E} x_{u} x_{v}+\sum_{v \in V} b_{v} x_{v}
$$

where $b \in\{0,1\}^{|V|}$
Find: $z \in \mathbb{F}_{2}^{|V|}$ such that $q(x)=2 z^{T} x, \forall x \in \mathcal{L}_{q}$.

## BACKGROUND

Quantum logic gates used in circuit $\mathcal{Q}_{N}$

$$
\begin{array}{ll}
\mathrm{H} & :=\left[\begin{array}{cc}
\frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} \\
\frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}}
\end{array}\right], \\
\mathrm{X}:=\left[\begin{array}{ll}
0 & \mathrm{~S} \\
1 & 0
\end{array}\right], \quad \mathrm{Y}:=\left[\begin{array}{cc}
1 & 0 \\
0 & -i
\end{array}\right], \\
\mathrm{Z}:=\left[\begin{array}{cc}
0 & -i \\
i & 0
\end{array}\right], \\
\left.\begin{array}{cc}
1 & 0 \\
0 & -1
\end{array}\right], \quad \mathrm{CZ}:=\left[\begin{array}{llll}
1 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 \\
0 & 0 & 1 & 0 \\
0 & 0 & 0 & -1
\end{array}\right]
\end{array}
$$

Definition: Let $G=(V, E)$ be a finite simple graph with $|V|=n$ and $|E|=m$. Suppose that a qubit is associated with each vertex of $G$. The $n$-qubit graph state of $G$ is

$$
\left|\phi_{G}\right\rangle:=\left(\prod_{(u, v) \in E} C Z_{u v}\right) H^{\otimes n}\left|0^{n}\right\rangle .
$$

## Example:



$$
\begin{aligned}
& V=\{1,2,3\}, E=\{(1,2),(2,3),(3,2),(2,1)\} \\
&\left|\phi_{G}\right\rangle=C Z_{12} C Z_{23} H^{\otimes 3}\left|0^{3}\right\rangle \\
&=\frac{1}{\sqrt{8}}(|000\rangle+|001\rangle+\ldots+|111\rangle) .
\end{aligned}
$$

Claim: For finite simple graph $G=(V, E)$, $\left|\phi_{G}\right\rangle$ is a stabilizer state for the stabilizer group generated by the operators

$$
g_{v}:=X_{v}\left(\prod_{(u, v) \in E} Z_{u}\right), \forall v \in V .
$$

Definition: Let $z \in\{0,1\}^{*}$. For a bit $z_{j}$ of $z$, define $m_{j}:=(-1)^{z_{j}}$. Describe $G$ and $L$ by


We define $L_{\text {odd }}$ and similarly
$L_{\text {even }}:=\{\ell \in L \mid \delta(\ell, u) \equiv 0(\bmod 2) \equiv \delta(\ell, v)\}$. Also define $m_{L}:=\prod_{j \in L_{\text {odd }}} m_{j}$.

## QUANTUM CIRCUIT FOR 2D-HLFP

Theorem: For every $N \geq 2$, there exists a quantum circuit $\mathcal{Q}_{N}$ of depth $d=O(1)$ which deterministically solves size $-N$ instances of 2D-HLFP.

## NONLOCALITY PROPERTY

Nonlocality: a form of correlation present in the measurement statistics of entangled quantum states that cannot be reproduced by local hidden variable models.
Example: Let $G$ describe the graph below


Let $b:=b_{u} b_{v} b_{w} \in\{0,1\}^{3}$ and define
$\mathcal{T}(b):=\left\{z \in\{0,1\}^{m} \mid\langle z| H^{\otimes m} S_{u}^{b_{u}} S_{v}^{b_{v}} S_{w}^{b_{w}}\left|\phi_{G}\right\rangle \neq 0\right\}$.
Claim: Let $b=b_{u} b_{v} b_{w} \in\{0,1\}^{3}$ and $z \in \mathcal{T}(b)$.
Then $m_{R} m_{B} m_{L}=1$. If $b_{u} \oplus b_{v} \oplus b_{w}=0$, then $i^{b_{u}+b_{v}+b_{w}} m_{u} m_{v} m_{w} m_{E} m_{R}^{b_{u}} m_{B}^{b_{v}} m_{L}^{b_{w}}=1$.
Lemma: The stabilizers of $\left|\phi_{G}\right\rangle$ are
$X_{u} X_{v} X_{w},-X_{u} Y_{v} Y_{w} X_{a} X_{c},-Y_{u} X_{v} Y_{w} X_{a} X_{b}$, and $-Y_{u} Y_{v} X_{w} X_{b} X_{c}$.

Measurements on classical local circuits cannot simultaneously satisfy these nonlocality identities!

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$\mathcal{Q}_{N}:$


