

# Shor's Algorithm and Its Impact on Modern Cryptography



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## Objectives

- To gain an intuitive understanding of how the quantum Fourier transform and Shor's algorithm work by visualizing the role of the roots of unity involved.
- To understand the impact of Shor's algorithm on the RSA cryptosystem, and thereby understand the importance of quantum computing in relation to modern cryptography.

## Introduction

The factoring problem is formalized in the following manner: given a composite odd integer  $N$ , find its prime factorization  $N = p_1^{\alpha_1} \cdot p_2^{\alpha_2} \cdot \dots \cdot p_n^{\alpha_n}$ . Since the number of operations required to find factors increases exponentially relative to the size of  $N$ , this is an infeasible problem. The fastest classical algorithm to date is the General Number Field Sieve (GNFS) algorithm, which has an asymptotic running time of  $O(e^{c(\log n)^{1/3}(\log \log n)^{2/3}})$  in terms of the length  $n$  of  $N$ . Because it is widely believed that  $P \neq NP$ , it is thought that no polynomial-time classical factoring algorithm exists. In quantum computing, however, this is possible: in 1995, Peter Shor formulated a quantum algorithm for factoring. One of the most commonly-used cryptosystems is RSA, which relies on the infeasibility of the factoring problem. In particular, the encryption and decryption function are defined modulo  $N$ , where  $N$  is of the form  $N = p \cdot q$ , for large primes  $p, q \in \mathbb{Z}_+$ . Because this cryptosystem plays a major role in the secure transmission of data, the potential ability to quickly factor  $N$  poses a threat. According to the National Institute of Standards and Technology (NIST), quantum computers will bring an end to modern cryptography as we know it.

## Quantum Computing Basics

*Qubits* - quantum bits, can be in any linear combination, or *superposition*, of the basis states  $|0\rangle$  and  $|1\rangle$ :  $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$ , where  $\alpha, \beta \in \mathbb{C}$  with normalization condition  $|\alpha|^2 + |\beta|^2 = 1$ .

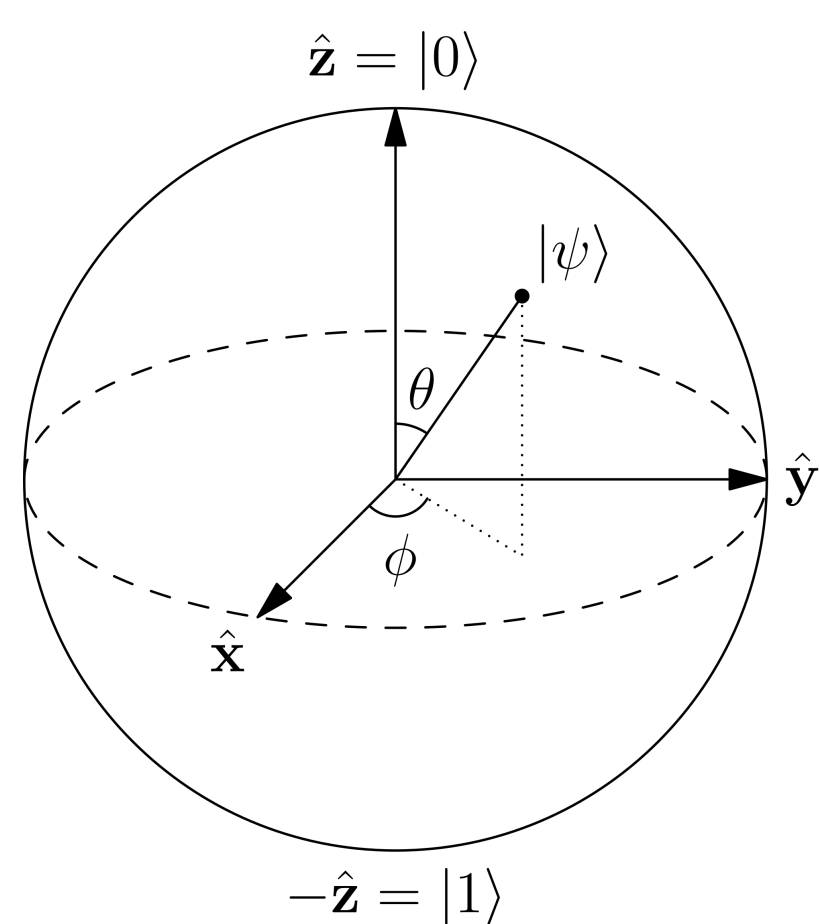


Figure 1: Bloch sphere

Quantum logic gates used in QFT:

*Hadamard gate*:  $H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$ .

*Controlled-phase gate*:  $R_k = \begin{bmatrix} 1 & 0 \\ 0 & e^{\frac{2\pi i}{2^k}} \end{bmatrix}$

## Shor's Algorithm

**I.** Reduction of factoring to order-finding algorithm:

*Input*: odd composite  $N \in \mathbb{Z}_+$

*Output*: non-trivial factors of  $N$

1. Choose a random  $a \in \mathbb{Z}_+$ ,  $a < N$ .
2. Compute  $\gcd(a, N)$  using Euclidean Algorithm.
3. If  $\gcd(a, N) \neq 1$ , return  $\gcd(a, N)$ .  
Else, use subroutine (II) to find the order  $r$ ,  $a^r \equiv 1 \pmod N$ .
4. If  $r$  odd or  $a^{r/2} \equiv -1 \pmod N$ , return to (1).  
Else, return  $\gcd(a^{r/2} + 1, N)$  and/or  $\gcd(a^{r/2} - 1, N)$ .

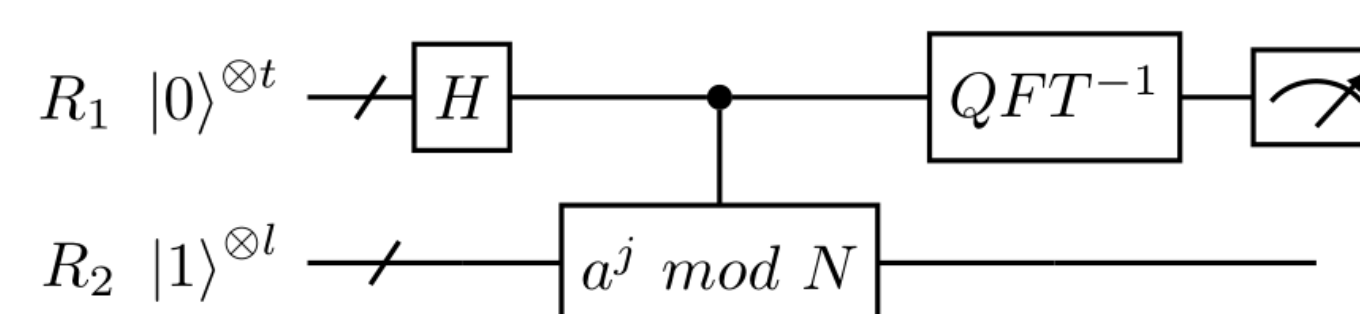


Figure 2: Quantum Subroutine Circuit

**II.** Quantum subroutine for order-finding:

*Inputs*:

- (i) black box transformation  $U_{a,N} : |j\rangle|k\rangle \rightarrow |j\rangle|a^j k \pmod N\rangle$  for  $a \in \mathbb{Z}_+$
- (ii)  $t$  qubits initialized to  $|0\rangle$ , where  $t := 2l + 1 + \lceil \log(2 + \frac{1}{2\epsilon}) \rceil$  and  $l := \lceil \log N \rceil$
- (iii)  $l$  qubits initialized to  $|1\rangle$

*Output*: order of  $a$  modulo  $N$

1. Apply the Hadamard gate to each qubit in  $R_1$ :  
 $H^{\otimes t}(|0\rangle^{\otimes t}) = (\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle))^{\otimes t} = \frac{1}{\sqrt{2}} \sum_{j=0}^{2^t-1} |j\rangle$ .
2. Apply  $U_{a,N}$  to each qubit in  $R_2$ :  
 $U_{a,N}(\frac{1}{\sqrt{2^t}} \sum_{j=0}^{2^t-1} |j\rangle|1\rangle) = \frac{1}{\sqrt{2^t}} \sum_{j=0}^{2^t-1} |j\rangle|a^j \pmod N\rangle$   
 $= \frac{1}{\sqrt{r 2^t}} \sum_{s=0}^{r-1} \sum_{j=0}^{2^t-1} e^{\frac{2\pi i j s}{r}} |j\rangle|u_s\rangle$   
 $=: |\psi\rangle$   
where  $|u_s\rangle = \frac{1}{\sqrt{r}} \sum_{j=0}^{r-1} e^{\frac{-2\pi i j s}{r}} |a^j \pmod N\rangle$  an eigenstate of  $U$  defined by  $U|x\rangle = |ax \pmod N\rangle$ .
3. Apply inverse QFT to  $R_1$ :

$$\begin{aligned} \text{QFT}^{-1}(|\psi\rangle) &= \frac{1}{\sqrt{r 2^t}} \sum_{s=0}^{r-1} \sum_{j=0}^{2^t-1} e^{\frac{2\pi i j s}{r}} \\ &\quad \left( \frac{1}{\sqrt{2^t}} \sum_{k=0}^{2^t-1} e^{\frac{-2\pi i j k}{2^t}} |k\rangle \right) |u_s\rangle \\ &= \frac{1}{\sqrt{r}} \sum_{s=0}^{r-1} \sum_{k=0}^{2^t-1} \alpha_{k,s} |k\rangle |u_s\rangle \end{aligned}$$

4. Measure  $R_2$  to choose an  $s$ , then measure  $R_1$  to obtain a value  $k$  for this  $s$ .
5. Apply continued fractions algorithm to  $\frac{k}{2^t}$  to find partial denominators  $r_0, r_1, \dots, r_l$ , and test  $r_i$  at each step to find the order  $r$ .

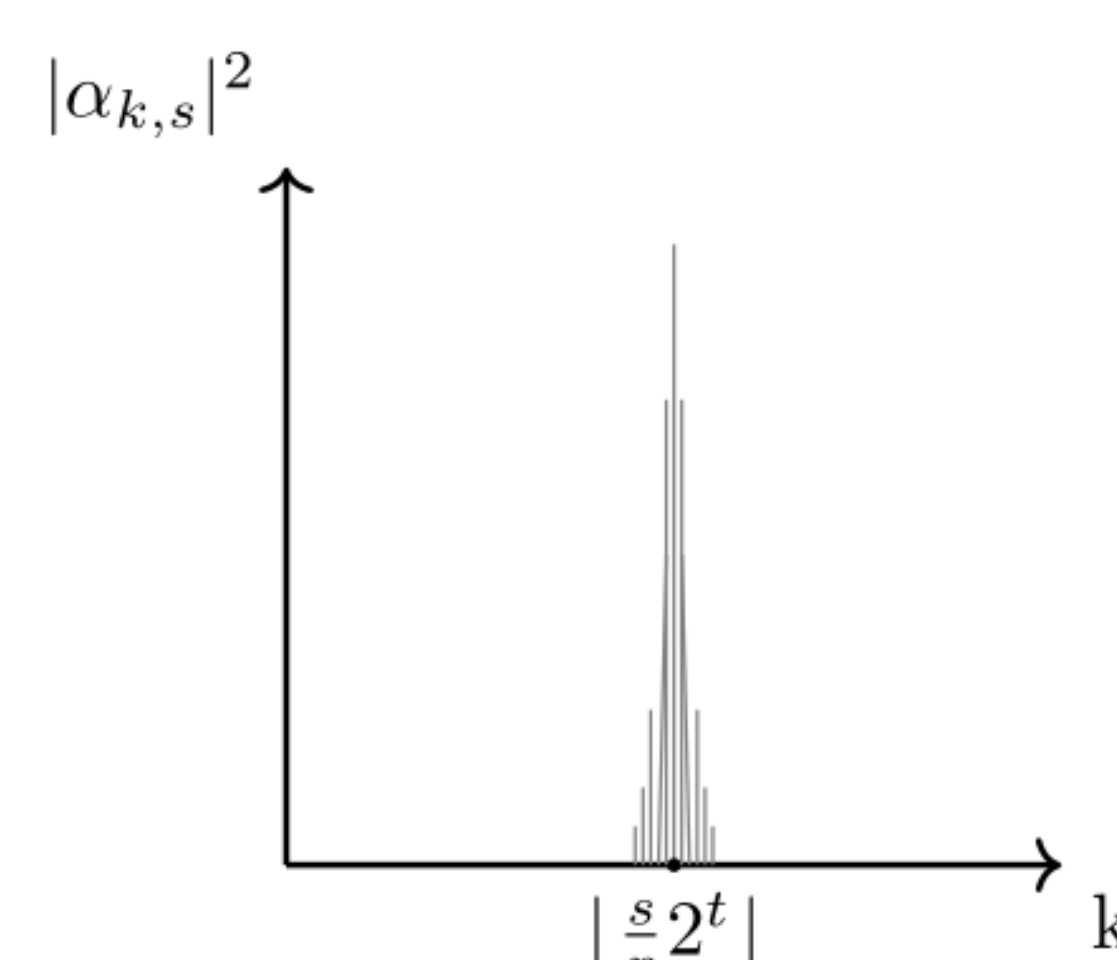


Figure 3: Probability Distribution for  $R_1$  Given  $R_2$

## Quantum Fourier Transform

An  $n$ -qubit system has basis states  $|0\rangle, \dots, |2^n - 1\rangle$ .

A state  $|j\rangle$  can be written in binary form

$$|j\rangle = j_1 \cdot 2^{n-1} + j_2 \cdot 2^{n-2} + \dots + j_n \cdot 2^0 = |j_1 j_2 \dots j_n\rangle.$$

The QFT acts on an input state  $|\psi\rangle$  by transforming each basis state  $|j\rangle$  by the following:

$$\begin{aligned} |j\rangle &\rightarrow \frac{1}{\sqrt{2^n}} \sum_{k=0}^{2^n-1} e^{\frac{2\pi i j k}{2^n}} |k\rangle \\ &= \frac{1}{\sqrt{2^n}} (|0\rangle + e^{2\pi i 0 \cdot j_n} |1\rangle) (|0\rangle + e^{2\pi i 0 \cdot j_{n-1}} |1\rangle) \dots \\ &\quad \dots (|0\rangle + e^{2\pi i 0 \cdot j_1} |1\rangle). \end{aligned}$$

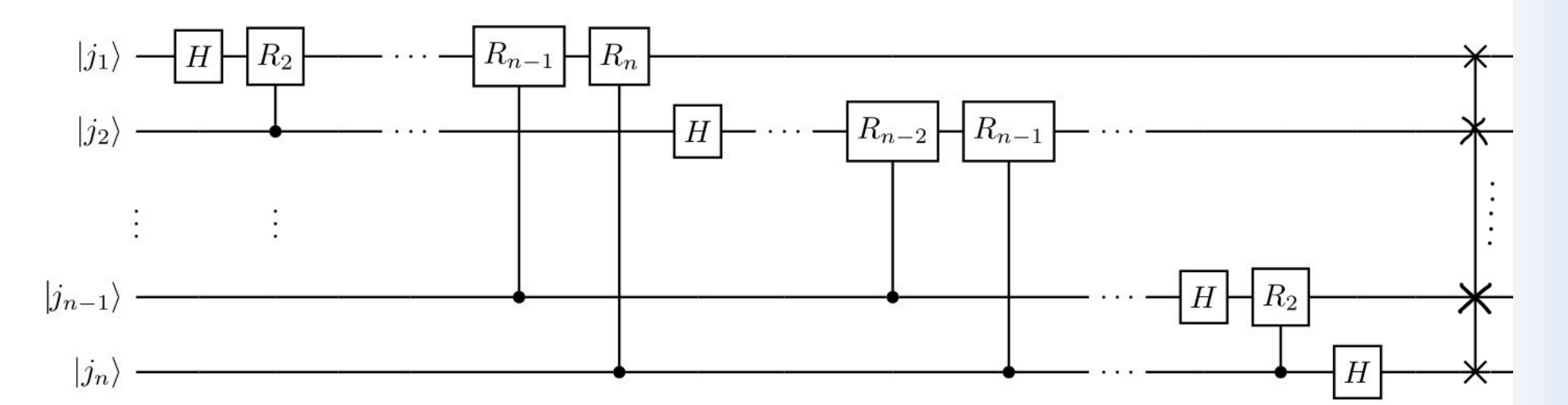


Figure 4: QFT Circuit Representation

## Roots of Unity

Let  $b + \delta := \frac{s 2^t}{r} - k$ , where  $b = \lfloor \frac{s 2^t}{r} - k \rfloor$ , and  $\omega := e^{\frac{2\pi i}{2^t}}$ . Then the amplitude in step II, 3. can be rewritten as

$$\alpha_{k,s} = \frac{1}{2^t} \sum_{j=0}^{2^t-1} \omega^{j(b+\delta)}.$$

Black roots correspond to  $b \geq 1$  and  $\delta = 0$ , shifted by a factor of  $\omega^b$  for  $b > 1$ . Grey roots correspond to error produced by  $\delta \neq 0$ , with shifting amount relative to exponent of  $\omega^j$ . For  $\delta = 0$ , the total sum is  $\alpha_{k,s} = 1$  resulting in no error, but is inexact otherwise.

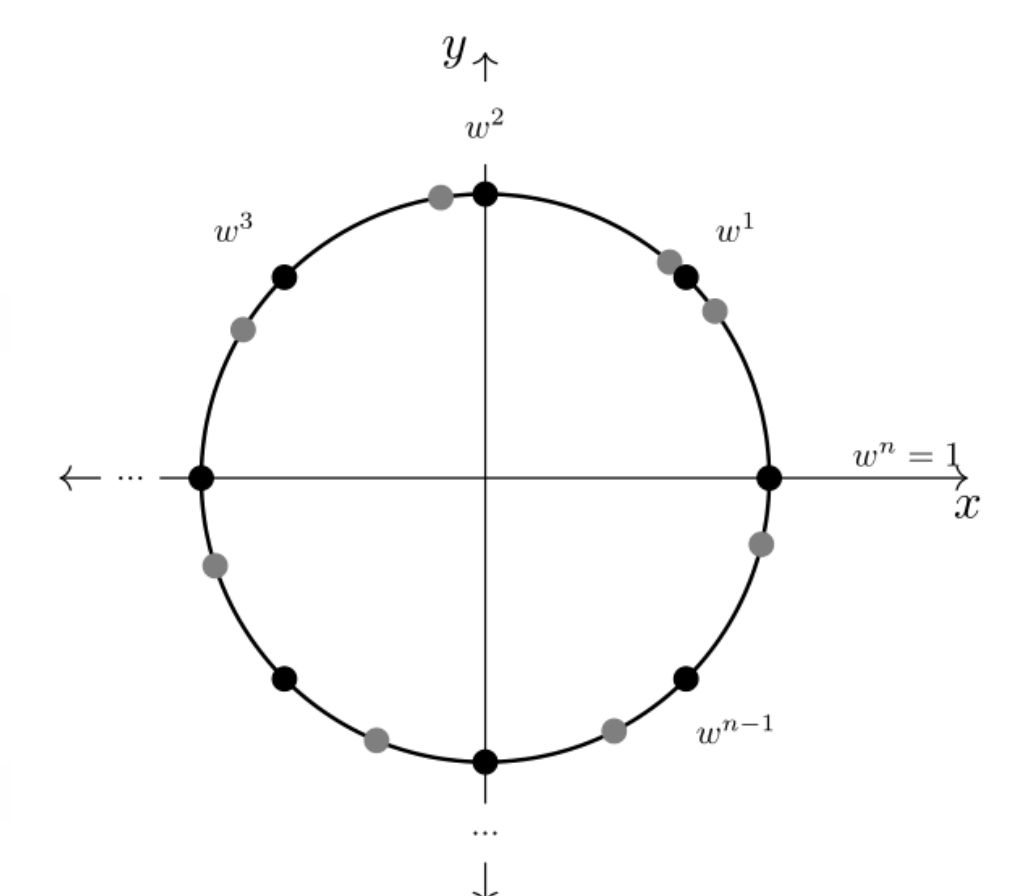


Figure 5: Roots of Unity for Quantum Subroutine

## Acknowledgements

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## References

- Nielsen, M.A. and Chuang I.L., *Quantum Computation and Quantum Information*, 2000
- Shor, Peter W., *Polynomial-Time Algorithms for Prime Factorization and Discrete Logarithms on a Quantum Computer*, 1995
- Weisstein, Eric W., *Number Field Sieve*, Wolfram MathWorld
- Rupert, Steven, et al., *Shor's Algorithm Simulator*, Colorado School of Mines, 2011