Lecture 2 – EM Waves Part I

EECS 598-002 Winter 2006 Nanophotonics and Nano-scale Fabrication P.C.Ku

Lecture 2-4 EM waves in nanoscale

- Interaction of EM waves with macroscopic media
 - Derivation of macroscopic Maxwell's equations
 - Boundary conditions
 - Basic properties of EM waves
- Examples of interaction between EM waves and macroscopic media
 - Interface phenomena
 - Evanescent waves
- Generation of EM waves
 - Antenna basics
 - Dipole antenna
 - Near field and far field
- More on the interaction between EM waves and macroscopic media
 - Diffraction and scattering
 - Wave guiding
- Scaling and symmetry in EM waves

Illustrations of polarization and magnetization

Polarization P



Magnetization M

Examples of different macroscopic media



Linear dielectric materials

- Isotropic media
 - Most of the materials can be regarded as isotropic dielectric materials. Another useful physical parameter refractive index *n* is defined as $n = \sqrt{\varepsilon}$
- Anisotropic media
 - For example, in calcite:

$$\mathbf{\varepsilon} = \begin{pmatrix} \varepsilon & 0 & 0 \\ 0 & \varepsilon & 0 \\ 0 & 0 & \varepsilon_z \end{pmatrix}$$



Because of two different refractive indices in different propagation directions, they are called "birefringent."

Magnetic materials

For examples, check http://www.irm.umn.edu/hg2m/hg2m_b/hg2m_b.html Diamagnetic $\mu < \mu_0$ Paramagnetic В $\mu > \mu_0$ Ferromagnetic ►H Antiferromagnetic hysteresis

Nonlinear optical (NLO) crystals



Frequency doubling through NLO.

Chiral materials



Boundary conditions



C and S are infinitesimal contour and surface.

$$\begin{split} \oint_C \vec{E} \cdot d\vec{l} &= -\frac{d}{dt} \int \vec{B} \cdot d\vec{S} = 0 \Rightarrow E_{1t} = E_{2t} \\ \oint_C \vec{H} \cdot d\vec{l} &= \int_S \left(\vec{J} + \frac{\partial \vec{D}}{\partial t} \right) \cdot d\vec{s} \Rightarrow \vec{n}_2 \times (\vec{H}_1 - \vec{H}_2) = \vec{J}_S \\ \oint_S \vec{D} \cdot d\vec{s} &= Q \Rightarrow D_{1n} - D_{2n} = \rho_S \\ \oint_S \vec{B} \cdot d\vec{s} &= 0 \Rightarrow B_{1n} - B_{2n} = 0 \end{split}$$



Applying boundary conditions at each interface, we have:

$$k_{1x} = k_{2x} = k_{3x}$$

In the case $n_1 > n_2$ and when $k_{1x} > \frac{2\pi n_2}{\lambda}$, $k_{2z} = \sqrt{\left(\frac{2\pi n_2}{\lambda}\right)^2 - k_{2x}^2}$ is imaginary.

Thin slab (cont.)

When the incident angle meets the total internal reflection condition at the 1-2 interface, there is no wave propagated in region 2 but there is still a finite transmission through the thin slab. The wave seems to "tunnel" through the slab!



For example, in the following plots, $n_1=1.5$ and $n_2=1$.

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In a general EM problem, we have different length scales:

 $\begin{array}{l} \lambda: \mbox{ wavelength} \\ r: \mbox{ size of the scatterer} \\ a: \mbox{ size of the diffracter} \\ R1, R2, \ldots: \mbox{ distances b/w two objects} \\ \end{array} \left\{ \begin{array}{l} \mbox{ Near: } R << \lambda \\ \mbox{ Intermediate: } R \sim \lambda \\ \mbox{ Far: } R >> \lambda \end{array} \right.$