Reconstruction from Digital Holograms by Statistical Methods

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Outline

- Background on holography
- Image model
- Conventional numerical reconstruction
- Statistical holographic reconstruction
- Simulation results
- Conclusions and future work

Conventional Holography

- "Record both amplitude and phase of a wave field"
- Hologram = recorded interference pattern between an object beam and a reference beam



Image Plane Holography



- Optical sectioning property as in confocal microscopy
- No xy scanning
- Larger source decreases coherence and improves spatial resolution

Challenges

- Finite source size limits spatial resolution
- Optical reconstruction inconvenient
- \Rightarrow Use statistical image reconstruction



Measurement Model



- Object beam, superposition: $u_0(\vec{r}') = \int h(\vec{r}',\vec{r})f(\vec{r}) d\vec{r}$
- Discretized object model $u_o(\vec{r}_i) = [\mathbf{A}\mathbf{x}]_i$, $a_{ij} = \int h(\vec{r}_i, \vec{r}) \chi_j(\vec{r}) d\vec{r}$
- Hologram intensity at recording plane

$$\begin{split} I(\vec{r}) &= |u_{\rm o}(\vec{r}) + u_{\rm ref}(\vec{r})|^2 \\ &= |u_{\rm o}(\vec{r})|^2 + |u_{\rm ref}(\vec{r})|^2 + u_{\rm o}(\vec{r})u_{\rm ref}^*(\vec{r}) + u_{\rm o}^*(\vec{r})u_{\rm ref}(\vec{r}) \end{split}$$

• Mean of *i*th measured sample

 $\begin{aligned} \mathsf{E}[Y_i] &= I(\vec{r}_i) + b_i \leftarrow \text{background dark current} \\ &= |u_o(\vec{r}_i) + u_{\text{ref}}(\vec{r}_i)|^2 + b_i = |[\mathbf{A}\mathbf{x}]_i + u_i|^2 + b_i, \qquad i = 1, \dots, N \end{aligned}$

Conventional Reconstruction Method

- Plane-wave assumption: $u_{ref}(\vec{r}) = U_{ref}e^{-i2\pi\vec{r}\cdot\vec{\alpha}}$
- Hologram intensity

$$I(\vec{r}) = |u_{\rm o}(\vec{r})|^2 + U_{\rm ref}^2 + U_{\rm ref}u_{\rm o}(\vec{r})\,{\rm e}^{\imath 2\pi\vec{r}\cdot\vec{\alpha}} + U_{\rm ref}u_{\rm o}^*(\vec{r})\,{\rm e}^{-\imath 2\pi\vec{r}\cdot\vec{\alpha}}$$

• Fourier transform of hologram intensity

$$I(\vec{\nu}) = I_o(\vec{\nu}) + U_{\text{ref}}^2 \delta(\vec{\nu}) + U_{\text{ref}} U_o(\vec{\nu} - \vec{\alpha}) + U_{\text{ref}} U_o^*(-\vec{\nu} - \vec{\alpha})$$



Statistical Model

Poisson model of the hologram measurement data:

 $Y_i \sim \text{Poisson}\left\{ \left| \left[\mathbf{A} \mathbf{x} \right]_i + u_i \right|^2 + b_i \right\}, \quad i = 1, \dots, N$

 $[\mathbf{A}\mathbf{x}]_i = \sum_{j=1}^P a_{ij} x_j$

- Known: imaging system A, data $\{Y_i\}$, reference beam $\{u_i\}$, dark current offset $\{b_i\}$
- Cost function: $\Psi(\mathbf{x}) = L(\mathbf{x}) + V(\mathbf{x})$
- Negative likelihood function: $L(\mathbf{x}) = \sum_{i=1}^{N} h_i([\mathbf{A}\mathbf{x}]_i)$

$$h_i(z) = (|z + u_i|^2 + b_i) - y_i \log(|z + u_i|^2 + b_i)$$

- Penalty: $V(\mathbf{x}) = \beta \sum_{i=1}^{r} \psi([\mathbf{C}\mathbf{x}]_i)$
- Penalized-likelihood (aka MAP) reconstruction:

$$\mathbf{x} \stackrel{\triangle}{=} rgmin_{\mathbf{x}} \Psi(\mathbf{x})$$

• No closed-form solution \Rightarrow need an iterative algorithm

Optimization Transfer Illustration



Optimization Transfer Principle

• S-step: find a majorizing surrogate function $\phi^{(n)}$

1.
$$\phi^{(n)}(\boldsymbol{x}^{(n)}) = \Psi(\boldsymbol{x}^{(n)})$$

2. $\phi^{(n)}(\boldsymbol{x}) \ge \Psi(\boldsymbol{x}), \quad \forall \boldsymbol{x} \in \mathbb{C}^{P}$

• M-step: minimize $\phi^{(n)}$ instead of Ψ

$$\boldsymbol{x}^{(n+1)} \stackrel{\triangle}{=} rgmin_{\boldsymbol{x}} \phi^{(n)}(\boldsymbol{x})$$

• Fact: algorithm monotonically decreases cost function (*cf.* Newton) $\Psi(\mathbf{x}^{(n+1)}) \leq \Psi(\mathbf{x}^{(n)})$



Log-likelihood Surrogate Functions



Quadratic Surrogate Function

S-step:

$$L(\mathbf{x}) = \sum_{i=1}^{N} h_i([\mathbf{A}\mathbf{x}]_i) \le \sum_{i=1}^{N} q_i([\mathbf{A}\mathbf{x}]_i; [\mathbf{A}\mathbf{x}^{(n)}]_i) \stackrel{\triangle}{=} Q(\mathbf{x}; \mathbf{x}^{(n)})$$

$$Q(\mathbf{x}; \mathbf{x}^{(n)}) \stackrel{\triangle}{=} \sum_{i=1}^{N} q_i([\mathbf{A}\mathbf{x}]_i; [\mathbf{A}\mathbf{x}^{(n)}]_i)$$

$$= L(\mathbf{x}^{(n)}) + \nabla L(\mathbf{x}^{(n)})(\mathbf{x} - \mathbf{x}^{(n)}) + \frac{1}{2}(\mathbf{x} - \mathbf{x}^{(n)})'\mathbf{A}'\mathbf{D}\mathbf{A}(\mathbf{x} - \mathbf{x}^{(n)})$$

$$\mathbf{D} \stackrel{\triangle}{=} \operatorname{diag}\left\{\breve{c}_i^{(n)}\right\}$$

M-step:

$$\boldsymbol{x}^{(n+1)} = \operatorname*{arg\,min}_{\boldsymbol{x} \in \mathbb{C}^{P}} Q(\boldsymbol{x}; \boldsymbol{x}^{(n)}) + V(\boldsymbol{x})$$

Surrogate *Q* is quadratic, so minimize by PCG algorithm.

Quadratic surrogates for regularizing penalty function also available. (*cf.* "half quadratic" methods)

Separable Quadratic Surrogate Algorithm

• Convexity of q_i allows separable surrogate (De Pierro, 1995):

$$q_{i}([\mathbf{A}\mathbf{x}]_{i}; z_{i}^{(n)}) = q_{i}\left(\sum_{j=1}^{P} \pi_{ij}\left[\frac{[a_{ij}(x_{j} - x_{j}^{(n)})]}{\pi_{ij}} + [\mathbf{A}\mathbf{x}^{(n)}]_{i}; z_{i}^{(n)}\right]\right)$$

$$\leq \sum_{j=1}^{P} \pi_{ij}q_{i}\left(\frac{[a_{ij}(x_{j} - x_{j}^{(n)})]}{\pi_{ij}} + [\mathbf{A}\mathbf{x}^{(n)}]_{i}; z_{i}^{(n)}\right)$$

 $\pi_{ij} \ge 0$ and $\sum_{j=1}^{P} \pi_{ij} = 1$

Para

for

• Separable quadratic surrogate function:

$$Q'(\mathbf{x}; \mathbf{x}^{(n)}) = \sum_{j=1}^{P} Q_j(x_j; \mathbf{x}^{(n)})$$
$$Q_j(x_j; \mathbf{x}^{(n)}) = \sum_{i=1}^{N} \pi_{ij} q_i \left(\frac{[a_{ij}(x_j - x_j^{(n)})]}{\pi_{ij}} + [\mathbf{A}\mathbf{x}^{(n)}]_i; z_i^{(n)} \right)$$
allelizable update: $x_j^{(n+1)} = \underset{x_j \in \mathbb{C}}{\operatorname{arg\,min}} Q_j(x_j; \mathbf{x}^{(n)}) + V_j(x_j; \mathbf{x}^{(n)}),$
$$i = 1 \qquad P$$

How Many Holograms?

- Number of measurement elements *N* is fixed by the recorder (*e.g.*, CCD camera pixels)
- Number of reconstructed pixels *P* is chosen by algorithm designer.
- Natural choice is P = N, but this is under-determined!

 ○(Measured data values are real, reconstructed field is complex.)
 ○Need P ≤ N/2 to avoid an under-determined problem.
 ○In conventional digital holography, FFT zero-padding is used.
- Possible strategies:
 - 1 recorded hologram; reconstruct half-size image: P = N/2
 - 1 recorded hologram; reconstruct full-size image: P = N (Must rely on regularization.)
 - 2 recorded holograms with different reference beams; reconstruct full-size holographic image P = N/2

Simulation Data: Complex Object

 $\times 10^4$

Original Object





Magnitude



Holograms









x 10⁴

Reconstruction Results: Complex Object



• Nonquadratic, edge-preserving regularization $\psi(t) = |t/\delta| - \log(1 + |t/\delta|)$, with δ chosen empirically

200 iterations of separable quadratic surrogate (SQS) algorithm

Contours of Marginal Objective Functions



 $h_i(z)$ for $z \in \mathbb{C}$, and $h_i([Ax^{(n)}]_i)$ vs n(Unregularized.)

Real-Valued Object (Constrained)



Hologram Data #2

1 Data Set, P=N/2(NMSE = 0.04)



2 Data Sets, P=N/2(NMSE = 0.03)

Summary

- Statistical method for reconstructing a complex-valued object field from real-valued hologram intensity data
- Poisson statistical model.
 Extendable to others, *e.g.*, Poisson+Gaussian (Snyder *et al.*)
- Optimization transfer monotonically decreases the cost function
- Preliminary simulations suggest improved image quality compared with conventional digital holography reconstruction method.
- Does not assume the reference beam is a plane-wave!
- But requires reasonably accurate reference beam model...
- Can incorporate constraints such as real-valued object field
- Accommodates "hologram reference beam diversity," reducing the problem of multiple minima

Future Work?

The usual story

- Use a space-variant model
- Regularization parameter selection
- Faster converging algorithms
- Test with real data

Specific to this topic

- Different types of digital holography (*e.g.*, Fresnel, Fourier)
- Phase retrieval problems