1131

Slab-Selective Spectral and Spectral-Spatial Prewinding RF Pulses

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Synopsis

We introduce a new type of small-tip angle prewinding RF pulse that compensates for spin dephasing attributed to off-resonance and is also slabselective, which can help limit the volume of coverage. We design purely spectral slab-selective pulses that prewind a limited global off-resonance bandwidth, and spectral-spatial slab-selective pulses that adapt the prewinding bandwidth spatially. We demonstrate these pulse designs in simulation and in experiments using a gel phantom with a distorted field and a volunteer's brain. Both pulses create sharp slab profiles, while the spectral-spatial pulse outperforms in terms of target magnetization phase.

Introduction

Spectral "prewinding" RF pulses provide spin-echo like re-phasing of spins in gradient echo sequences¹. A small-tip angle² purely spectral prewinding pulse recovered a limited bandwidth of off-resonance for imaging in the small-tip fast recovery sequence (STFR), a steady state sequence with similar contrast to balanced-SSFP³. A spectral-spatial pulse increased the effective prewinding bandwidth by varying the bandwidth spatially using a 2D field map⁴. Here, we demonstrate for the first time the incorporation of spectral prewinding into slab-selective pulses; these are practically important for limiting the FOV in the slab dimension.

Methods

We designed slab-selective prewinding pulses to excite target pattern, **d**, by sampling the function:

$$d(z, f) = p(z)e^{2\pi i \operatorname{TE} f}$$

where z is the slab dimension, f is a range of off-resonance frequencies within a 3D field map, p(z) is the slab-select envelope, and echo time TE is half the steady-state free precession time for STFR. As written, **d** defines the purely spectral³ slab-selective pulse, and we replicate it over transverse dimensions x and y to create a 4D spectral-spatial⁴ slab-selective pulse.

We designed the RF pulses by solving a small-tip angle approximation (STA)² constrained optimization problem:

$$\hat{\mathbf{b}} = \arg\min_{\mathbf{b}} ||\mathbf{A}\mathbf{b} - \mathbf{d}||_{\mathbf{W}}^2 \text{ s. t. } ||\mathbf{b}||_{\infty} \le b_{\max}$$

Where **A** is the STA system matrix, **W** is a weighting matrix, and b_{max} is the peak RF amplitude limit (0.2 Gauss). For the purely spectral slab-selective pulse, **W** is unity except for the transition regions between in-slab and out-of-slab areas as well as for frequencies outside of the field map within the slab which are allowed to vary freely. For the spectral-spatial version **W** is a 4D weighting function, extending ref. (4) to 3D. Figure 1 shows a 3D field map and an example of a 2D spectral target pattern and weighting matrix.

We obtain slab selectivity by applying excitation gradients during the RF pulse that sweep through kz-space. In the spectral-spatial slab-selective pulse, we employ additional x and y gradients that span the transverse excitation k-space as kx-ky blips, making a 3D spokes trajectory. We explored multiple spokes trajectories for the RF designs to balance between spectral and slab selectivity, while still maintaining a short pulse length. We generally found that four spokes at $\pm k_{x,max}$ and $\pm k_{y,max}$, repeated twice, worked sufficiently well--however spoke selection remains an open area of research.

To mimic realistic field distortions, we attached a small piece of metal to a gel phantom, and then acquired a field map. We then designed 5.04ms spectral and spectral-spatial slab-selective pulses to excite a 4cm slab. For both pulse types, we designed tip-down and tip-up pulses for the STFR sequence, where the tip-up pulse is formed by negating and reversing a pulse designed to excite the simulated magnetization magnitude and phase accrued at the end of free precession⁵. We used these pulses to acquire STFR images [FA/TE/TR=16°/3.6ms/18.3ms, FOV=24cm/24cm/24cm, matrix size 256x256x60] using a birdcage single channel T/R head coil on a 3T GE MR750 scanner. We repeated this experiment in a volunteer's brain. The 4D spectral-spatial slab-selective pulses were the most computationally demanding, taking about 3 min to compute online.

Results

Figure 2 shows the tip-down RF and gradient waveforms and the 3D spokes k-space trajectory for the spectral-spatial pulse used in the phantom. Figure 3 shows the simulated magnetization and phase for the target design pattern and the spectral and spectral-spatial slab-selective pulses in the phantom at TE after a single excitation (i.e., not in steady-state). The simulated full-FOV_z magnitude NRMSE and in-slab phase RMSE were 0.31/15.5° and 0.15/2.1° for the spectral and spectral-spatial pulses, respectively. Figure 4 shows the experimental images in the phantom.

Figure 5 shows the experimental images from the analogous experiment repeated in a volunteer. The magnitude NRMSE and phase RMSE computed from simulated images (not shown) were 0.25/7.2° and 0.12/1.4° for the spectral and spectral-spatial pulses, respectively.

Discussion and Conclusion

Prewinding slab-selective pulses are challenging because there is an inherent trade-off between slab-selectivity and phase prewinding. In simulation, spectral-spatial pulses balanced these competing needs more effectively than purely spectral slab-selective pulses by achieving lower magnitude NRMSE and phase RMSE. The spectral pulse achieved a sharper slab profile in the phantom than the spectral-spatial pulse, but again was outperformed by the

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spectral-spatial pulse slab profile in the human experiment as well as all phase images (phantom and human). Performance differences between simulation and experiment in the phantom might be attributed to steady-state behavior of the STFR sequence, as the simulation is conducted at TE directly after a single tip-down pulse.

Future work will involve optimizing the 3D excitation trajectory to balance slab-selectivity and spectral prewinding. Additionally, strategies to improve the tip-up pulse design could better match the simulated magnetization with the experimental steady-state results using STFR.

Acknowledgements

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References

1. Assländer J, Glaser SJ, Hennig J. Spin echoes in the regime of weak dephasing. Magn Res Med 2016;75:150–160.

2. Pauly J, Nishimura D, Macovski A. A k-space analysis of small-tip angle excitation. J Magn Reson 1989;81:43–56.

3. Sun H, Fessler JA, Noll DC, Nielsen JF. Balanced SSFP-like steady-state imaging using small-tip fast recovery (STFR) sequence with a spectral pre-winding pulse. Magn Res Med 2016;75:839–844.

4. Williams SN, Nielsen JF, Fessler JA, Noll DC. Design of spectral-spatial phase prewinding pulses and their use in small-tip fast recovery steady-state imaging. Magn Res Med 2017; DOI 10.1002/mrm.26794.

5. Nielsen JF, Yoon D, Noll DC. Small-tip fast recovery imaging using non-slice-selective tailored tip-up pulses and radiofrequency-spoiling. Magn Res Med 2013;69:657–666.



Figure 1. Full and within-slab 3D field maps of gel phantom with small piece of metal (top row) and 2D target pattern and weighting matrix (bottom row) for a spectral slab-selective pulse. On the bottom row, the horizontal red lines highlight the bandwidth of off-resonance frequencies found within the slab, although the entire object field map covers a larger range. The weighting matrix contains "don't care" regions for frequencies outside of this range within the slab, as well as "don't care" transition regions between in in-slab and out-of-slab in the spatial dimension, z.



Figure 2. RF pulse magnitude (top left), phase (middle left), excitation gradients (bottom left), and excitation k-space in 3D space (right) for the spectralspatial slab-selective pulse used to image the phantom. The RF pulse is peak-amplitude constrained to 0.2 Gauss (red line on top left). The excitation trajectory is a spokes trajectory where the four side spokes are located at $\pm k_{x,max}$ and $\pm k_{y,max}$. The entire spokes trajectory is repeated twice during a single excitation.



Figure 3. Target design pattern and Bloch simulated magnetization magnitude (top row) and phase (bottom row) at the echo time after one tip-down pulse displayed at the center slice in *xy*, *xz*, and *yz* planes in the phantom with a small piece of metal. The target magnitude is uniform within the slab and the target phase is zero, to provide spin-echo like contrast in the STFR sequence.



Figure 4. Experimental magnetization magnitude (top left row) and phase (bottom left row) for the steady-state STFR images in the *xy*, *xz*, and *yz* planes in the gel phantom with a small piece of metal. On the right are the *x* and *y* slab profiles for both spectral (dotted line) and spectral-spatial (solid line) pulses. The spectral slab-selective pulse has the best slab profile while the spectral-spatial slab-selective pulse has the flattest phase image.



Figure 5. Experimental magnetization magnitude (top row) and phase (bottom row) for the steady-state STFR images in the *xy*, *xz*, and *yz* planes in the human brain. On the right are the *x* and *y* slab profiles for both spectral (dotted line) and spectral-spatial (solid line) pulses. Here, the spectral-spatial slab-selective pulse has the best slab profile and the flattest phase.

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